

# International Advanced School on Muon Spectroscopy

## Stimulation Methods 2: Laser

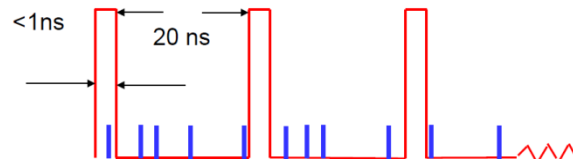
Koji Yokoyama

ISIS Neutron and Muon Source, STFC

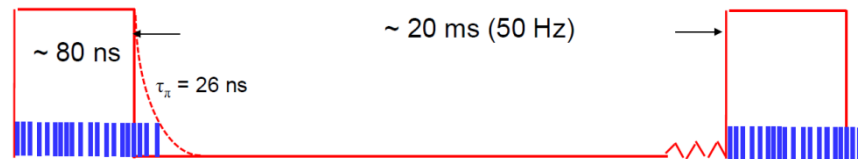
# How to use pulsed stimuli?

Muon sources come in two 'flavours' ...

- 'continuous beam' (PSI, TRIUMF), muons arrive randomly



- 'pulsed beam' (ISIS, J-PARC), all muons in short bunches



# Pulsed stimulations in $\mu$ SR...

Many possible, but these are typical ...

- Radio-frequency (RF) resonance
- Currents
- Magnetic fields
- Electric fields
- Acoustic resonance
- Illumination (flash lamp and laser)
- ...



Photoexcited  $\mu$ SR ... "photo- $\mu$ SR"

# photo- $\mu$ SR: What is it ?

- Light-pump muon-probe experiment
  - utilising standard  $\mu$ SR techniques
  - muon probes the photo-induced effect
  - track dynamics
- Has been done in various muon facilities
- Still in growing phase

## Representative works:

### KEK

- S. Chu, et al. Phys. Rev. Lett. 60, 101 (1988).
- K. Shimomura, et al. Hyperfine Interactions 120–121, 595 (1999).

### TRIUMF

- I. Fan, et al. Phys. Rev. B 77, 035203 (2008).

### PSI

- T. Prokscha, et al. Scientific Reports 3, (2013).

### RIKEN-RAL

- R. Kadono, et al. Phys. Rev. Lett. 73, 2724 (1994).

### ISIS

- K. Wang, Nat Mater 16, 467 (2017).
- K. Yokoyama, Phys. Rev. Lett. 119, 226601 (2017).

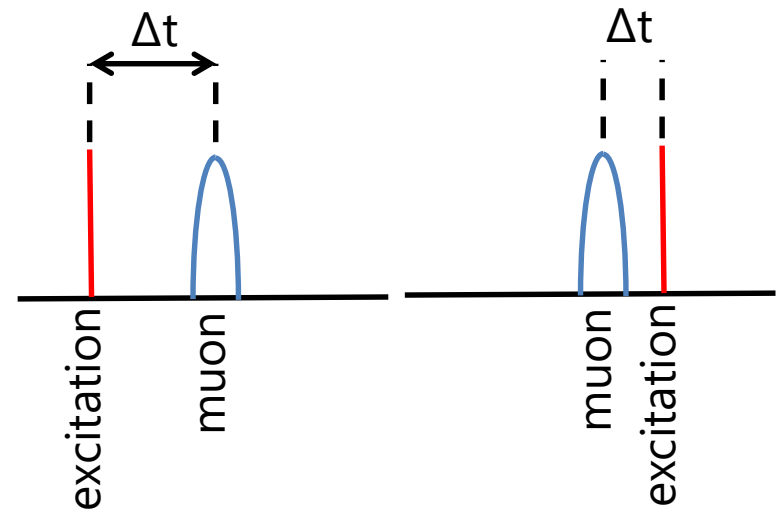
# Outline of talk

- Techniques
  - How photo- $\mu$ SR technique works
  - Light sources
  - Things to remember for designing experiments
- Case studies
  - Carrier kinetics measurements in semiconductor
  - Excited molecules 1: organic semiconductor
  - Excited molecules 2: reaction kinetics

# Pulsed muon + pulsed light source

Advantages:

- Arbitrary pulse timing:  
wide range of dynamics



- Pulsed laser ... lots of photons packed in a short pulse
- Smaller disturbance on the system
  - less local heating
  - Can provide large stimulation

# Light source

## Pulsed

Lasers



+ Flashlamps

## Continuous

Lamps

LED

+ CW laser

## Photoexcitation by pulsed laser light:

- High peak intensity
- Short pulse
- Wavelength conversion
  - optical parametric oscillator, dye laser, etc.
- Monochromatic
  - excite specific energy
  - polarisation control (e.g. circularly polarized light)
- Collimated beam
  - easy to calculate fluence

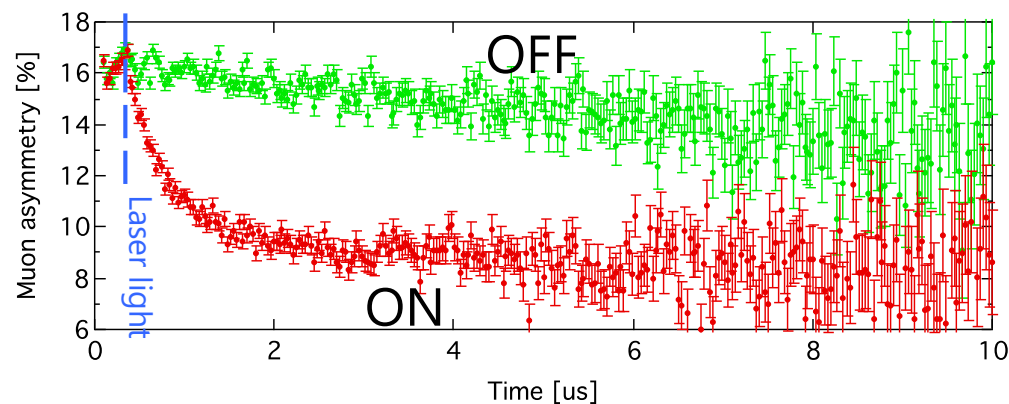
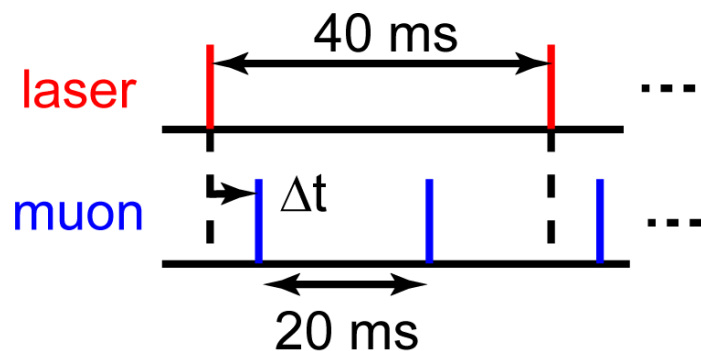
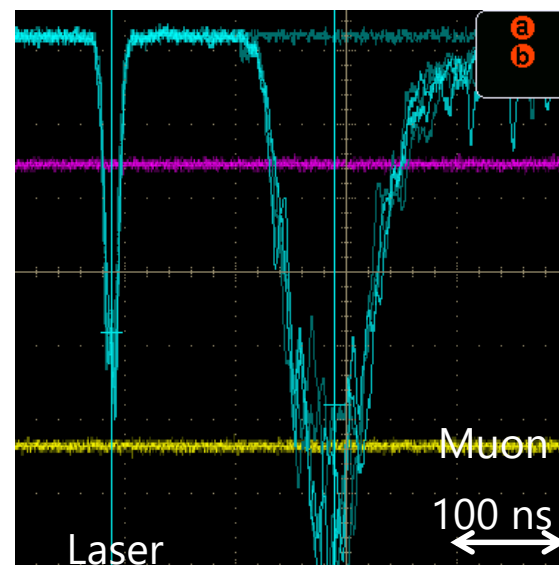
... enables more quantitative measurements

## Disadvantage:

- Health & Safety
- Complicated optical setup
- Expensive

# Nanosecond lasers match up nicely

- Time scales. Pulse FWHM:
  - muon 70 ns, laser 10 ns
  - and data bin <16 ns
- Repetition rate:
  - muon 50 Hz, laser 25 Hz

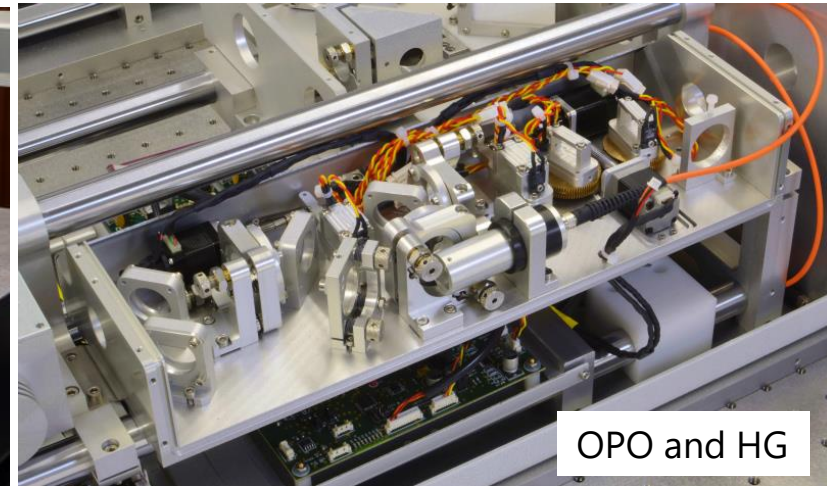
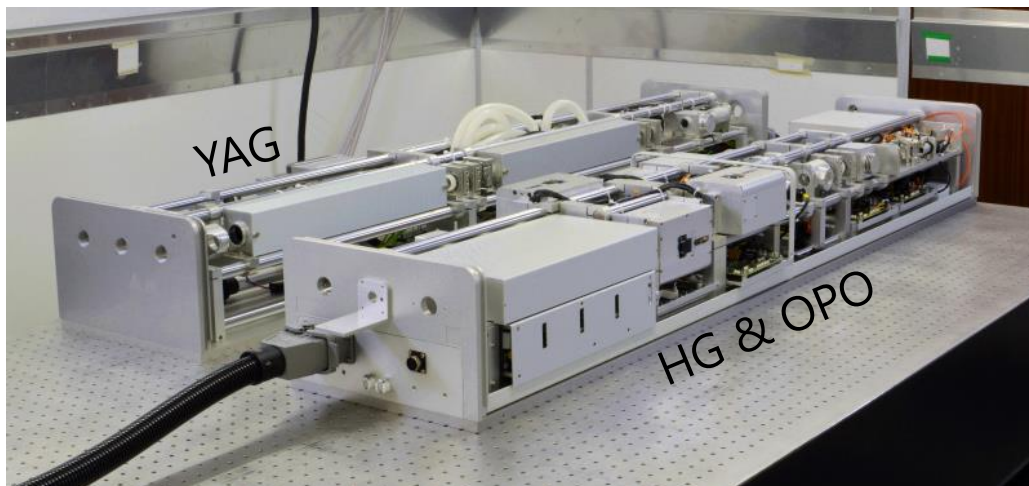
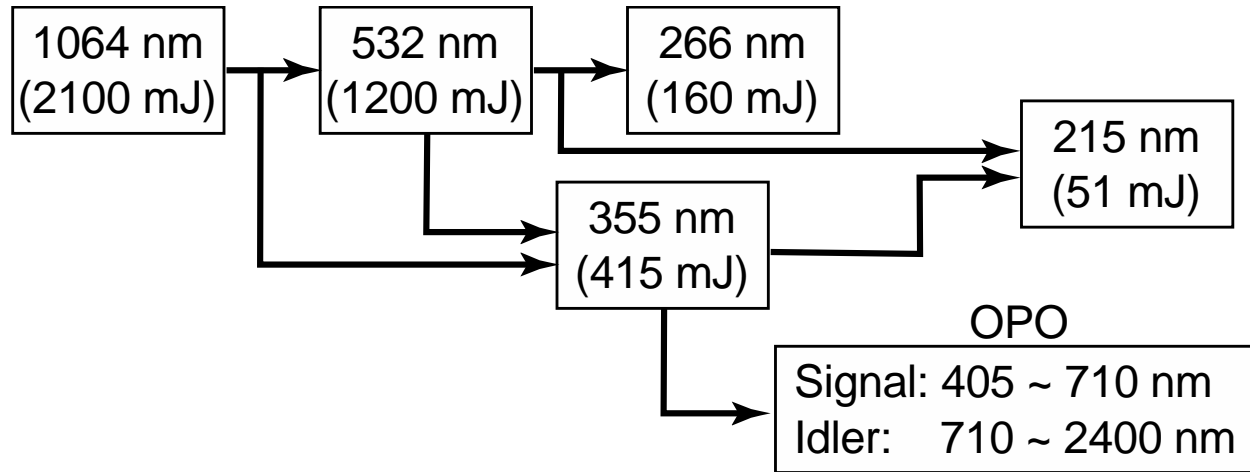


- Measure Light-ON/OFF spectrum alternately
- Can remove long-term drift

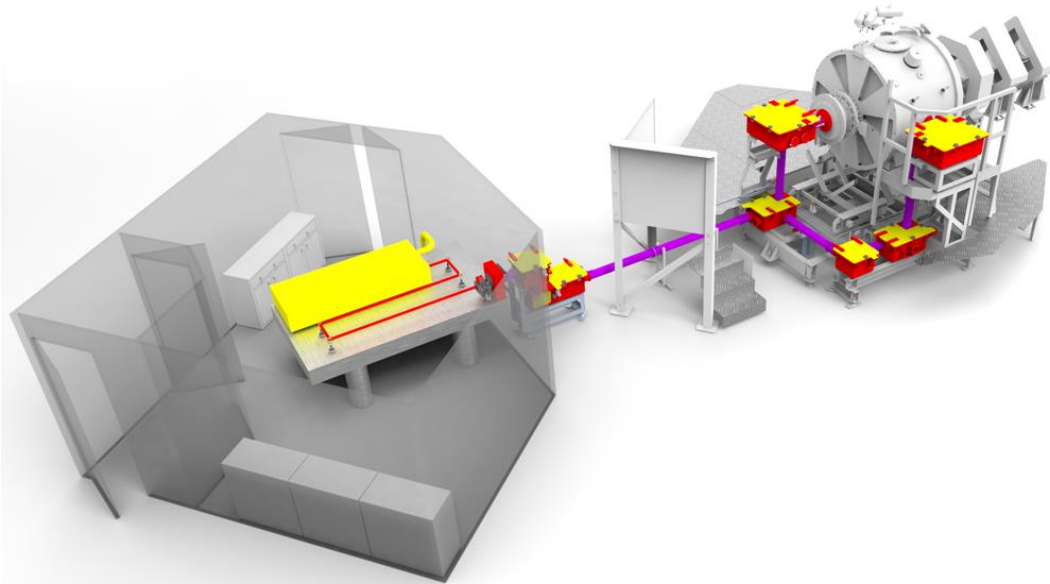


# Our laser system

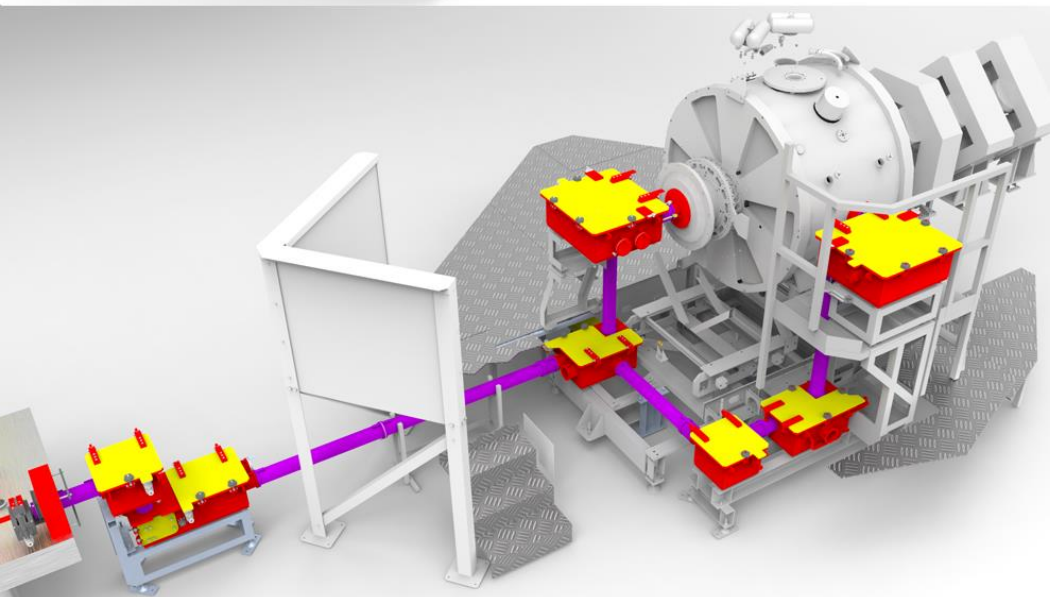
ns Q-switched Nd:YAG & harmonic generation units  
Optical parametric oscillator (OPO)



# HiFi Laser System

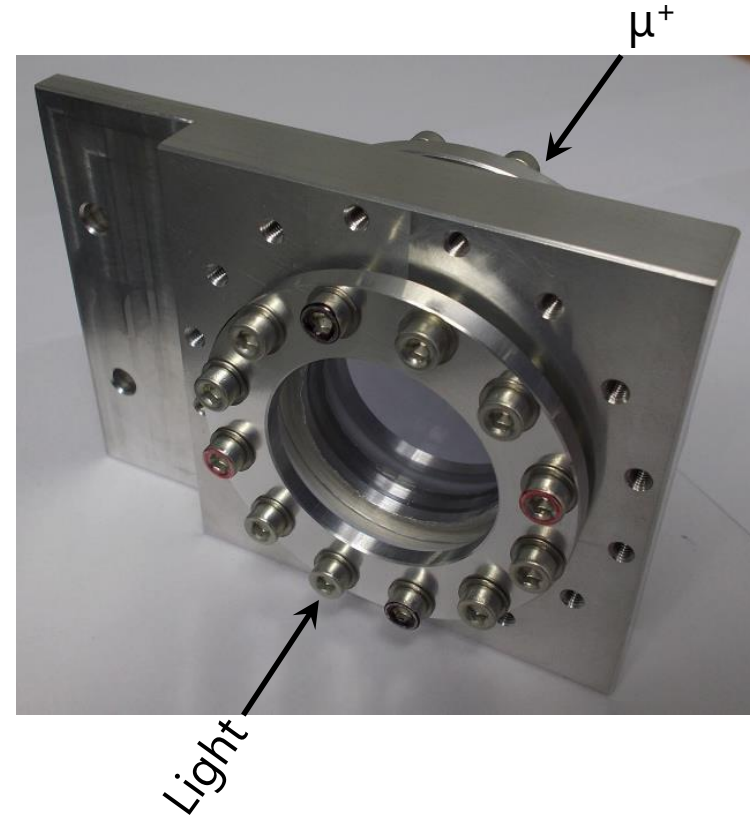
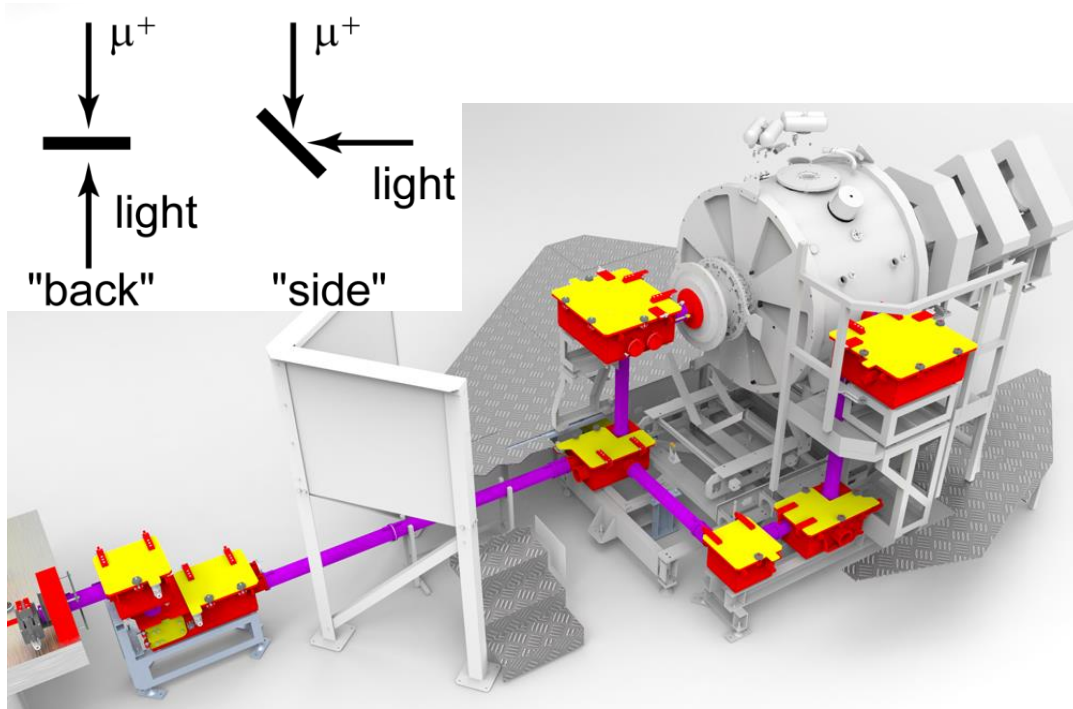


- Two optical setups in ISIS
  - ARGUS ... good for low field measurements
  - HIFI ... high-field, easier optics handling



- Laser cabin in the down stream of HiFi
- Laser beam transported through tubes
- Beam Entry Chamber attached for photo- $\mu$ SR experiments

# Geometry

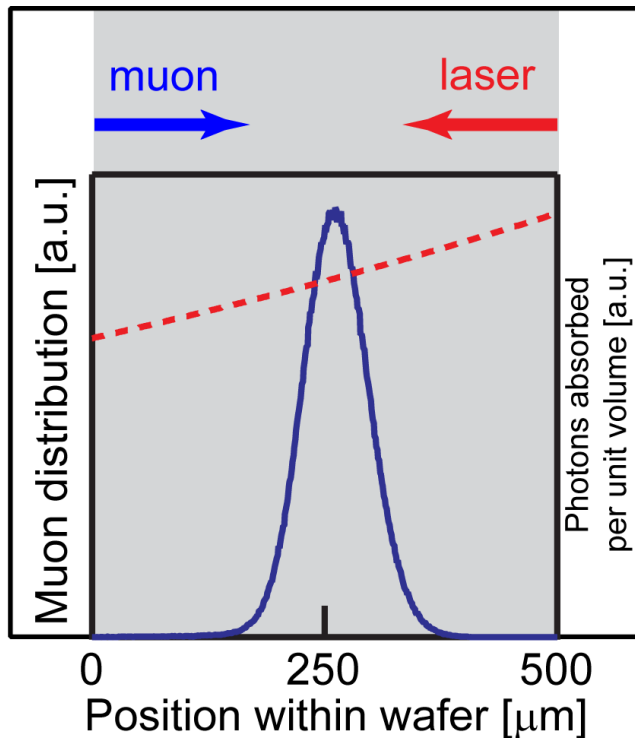


- Two possible beam geometries
- We normally use the "back-pump" geometry

- Sample cell with optical window
- Mounted on CCR

# Spatial overlap

Muon & photons: need spatial overlap in sample



- Muon: Gaussian-like distribution
  - can be calculated with *e.g.* musrSim
- Photons: (if optically uniform) exponential decay ( $e^{-\alpha z}$ ) with absorption coefficient  $\alpha$  [ $\text{cm}^{-1}$ ]
  - need  $\alpha$  for getting fluence at muon position
  - $\alpha$  depends on wavelength
- Lateral size:
  - muon beam size:  $\sim 10$  mm
  - laser: variable (better expanded for entire coverage)

# Before designing a photo- $\mu$ SR experiment...

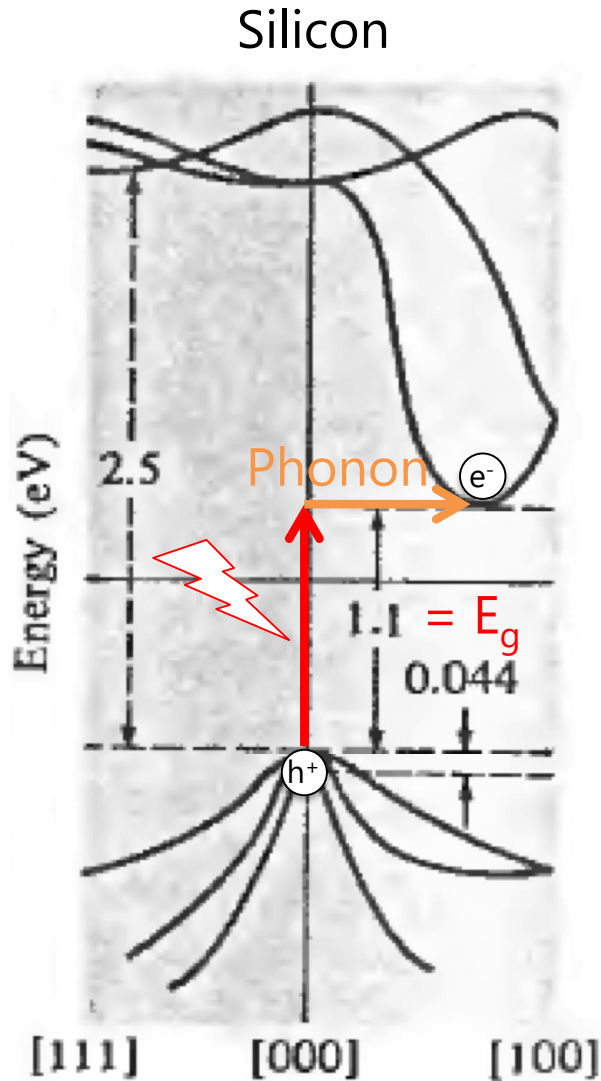
- Optically uniform sample, if possible
- How likely a muon can encounter the photo-induced effect?
  - *e.g.* injected excess carriers in semiconductors are v. mobile ( $e^-$  in Si can diffuse 1 mm in 1 ms)
- Absorption length? And how much laser power do you need?

## Case study #1

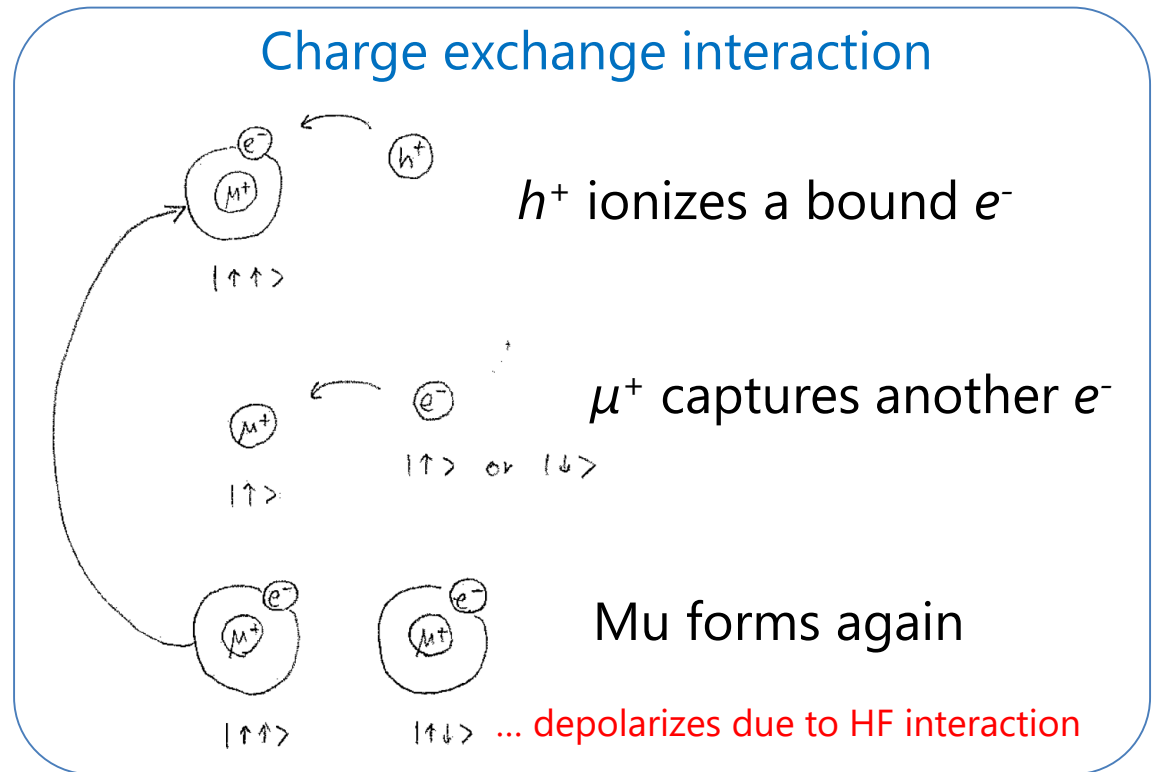
# Measuring excess carrier kinetics in semiconductors (Si & Ge)

- K. Yokoyama, J. S. Lord, J. Miao, P. Murahari, and A. J. Drew, Phys. Rev. Lett. 119, 226601 (2017).
- K. Yokoyama, J. S. Lord, P. W. Mengyan, M. R. Goeks, and R. L. Lichti, ArXiv:1906.07464 [Cond-Mat] (2019).

# Excess carriers depolarises $\mu^+$ spin



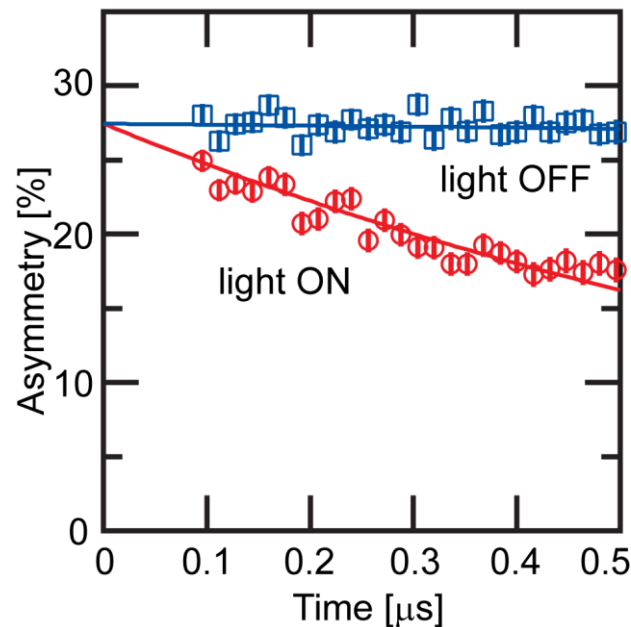
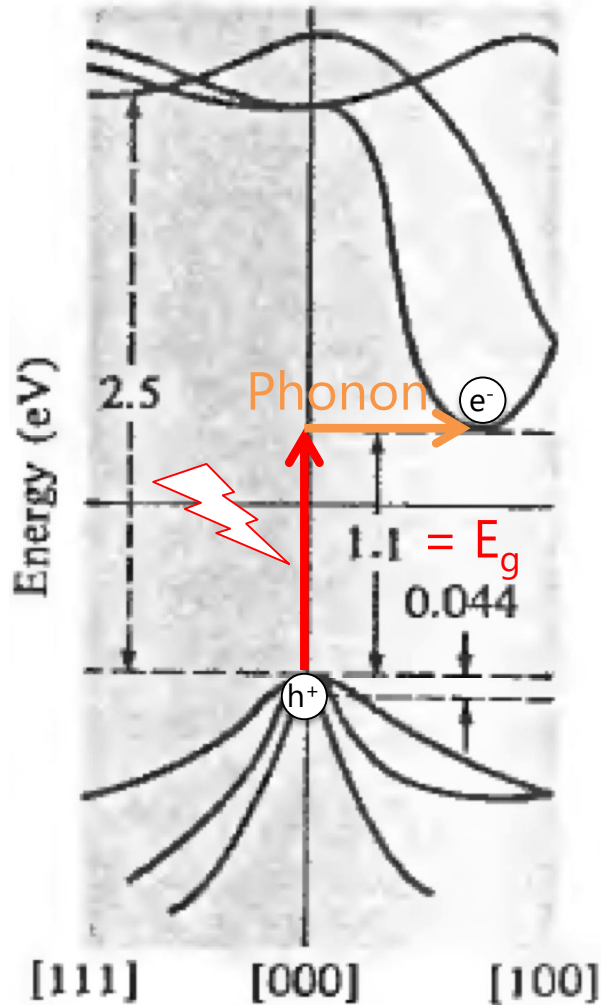
- Photon creates  $e-h$  pair (excess carriers)
- Indirect transition (requires phonon)



- Carrier injection causes the muon spin depolarisation

# Excess carriers depolarises $\mu^+$ spin

- Indirect transition (requires phonon)
- Photon creates  $e-h$  pair (excess carriers)

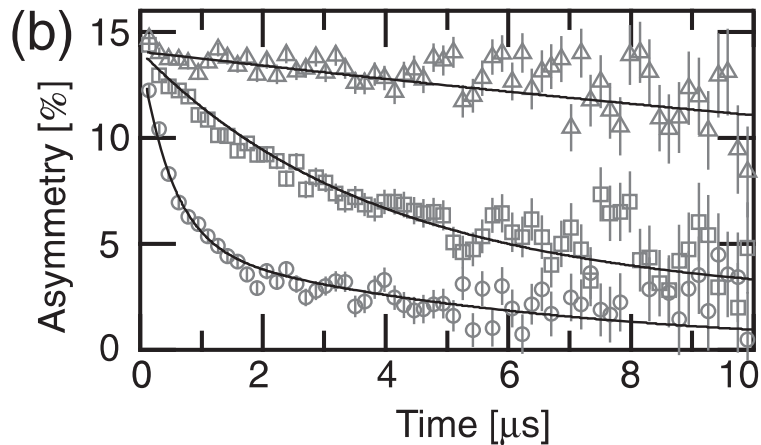


light pulse arrives just before muon

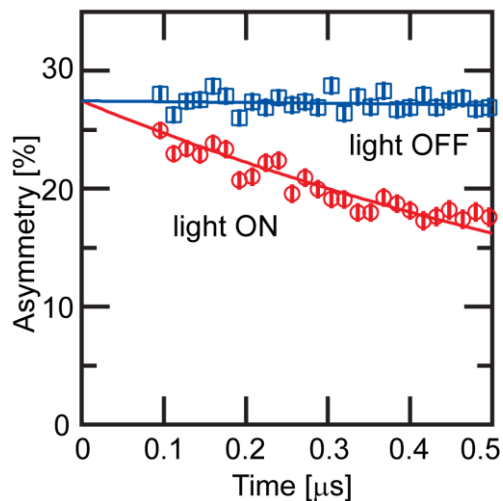
- Carrier injection causes the muon spin depolarisation



# Characterise $\Delta n$ with relaxation rate

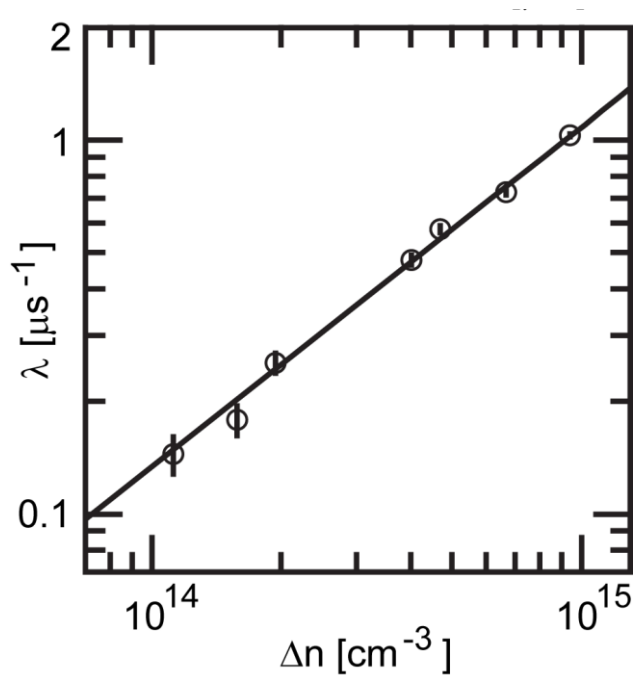


- More carrier injection gives faster relaxation
  - ... can we use this to measure excess carrier density  $\Delta n$  [ $\text{cm}^{-3}$ ]?
  - ... Relaxation rate  $\lambda$  for light ON spectrum is a measure of  $\Delta n$



- Need: relationship between  $\lambda$  vs.  $\Delta n$
- $\Delta n(\Delta t=0)$  is measurable
  - Size of beam spot
  - Pulse energy
  - Absorption coefficient
- Can get  $\lambda$  vs.  $\Delta n$  by changing pulse energy
- Assumption:  $\Delta n$  stays const. within the time window (0.5  $\mu\text{s}$ )

# Making a calibration curve

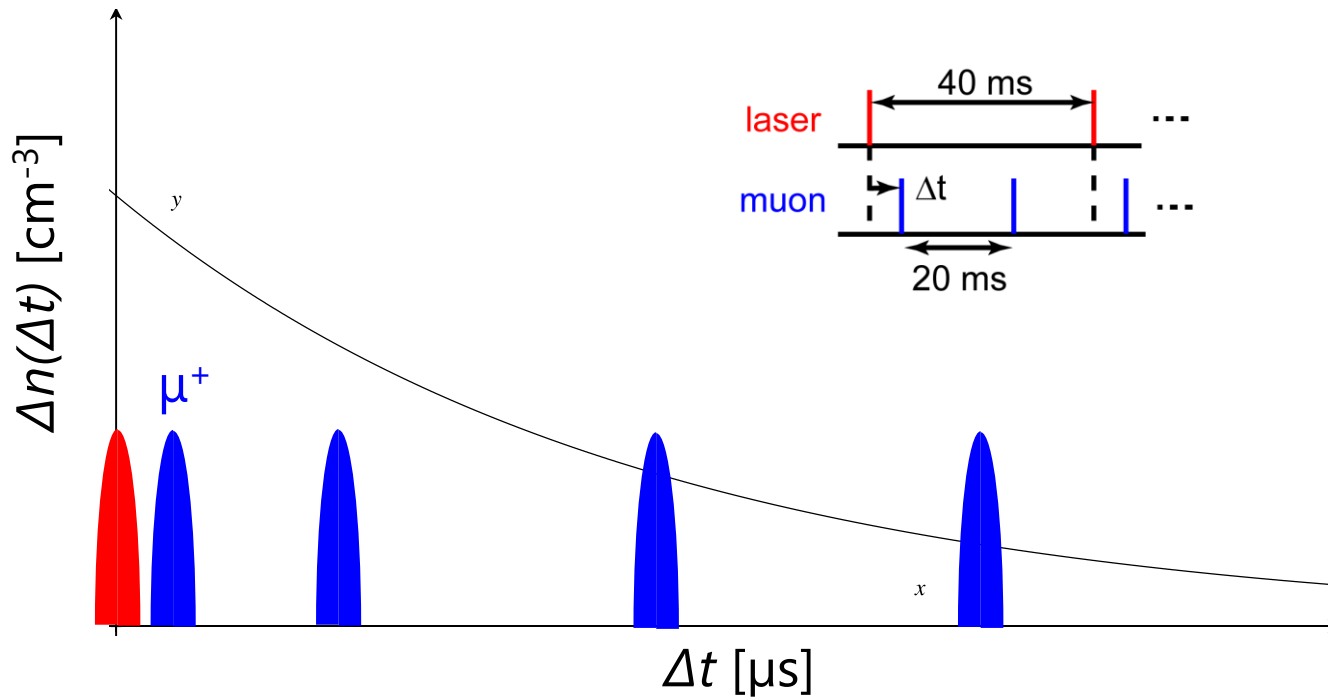


- $\lambda$  vs.  $\Delta n$ : data points can be fit with a power law:

$$\lambda = \beta \left( \frac{\Delta n}{\Delta n_0} \right)^\alpha$$

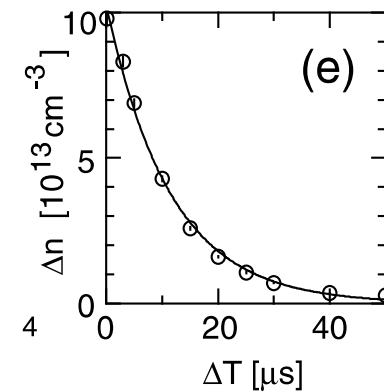
- The curve allows us to calculate  $\Delta n$  from a measured  $\lambda$

# Carrier lifetime spectrum



Now fix  $\Delta n(\Delta t=0)$  and change the delay time  $\Delta t$ , measure  $\lambda$ , and get  $\Delta n$

... gives a carrier lifetime spectrum

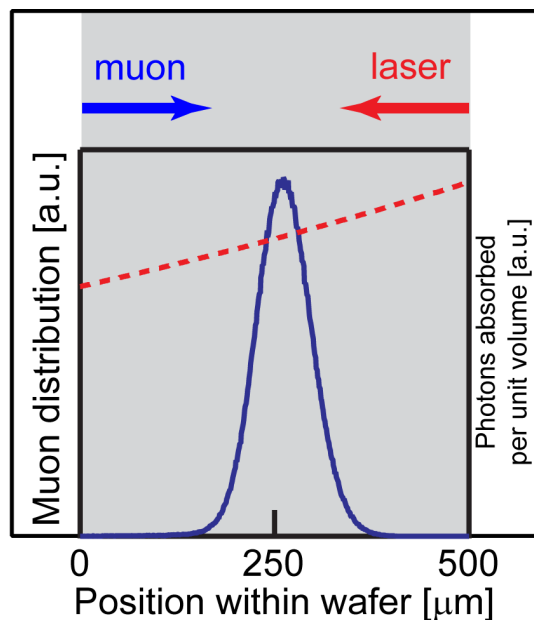


# Carrier lifetime spectrum

Yet another technique for lifetime spectroscopy  
*e.g.* photoconductance, microwave reflection etc.

- Mu: common in semiconductors  
... applicable to other systems?
- Something new in photo- $\mu$ SR ?

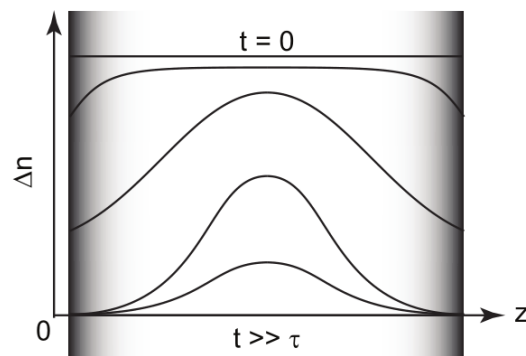
# Muons are localised



- Spatially localized probe
- Lifetime spectrum measured at the centre
- Can get transport properties?

... to test this,

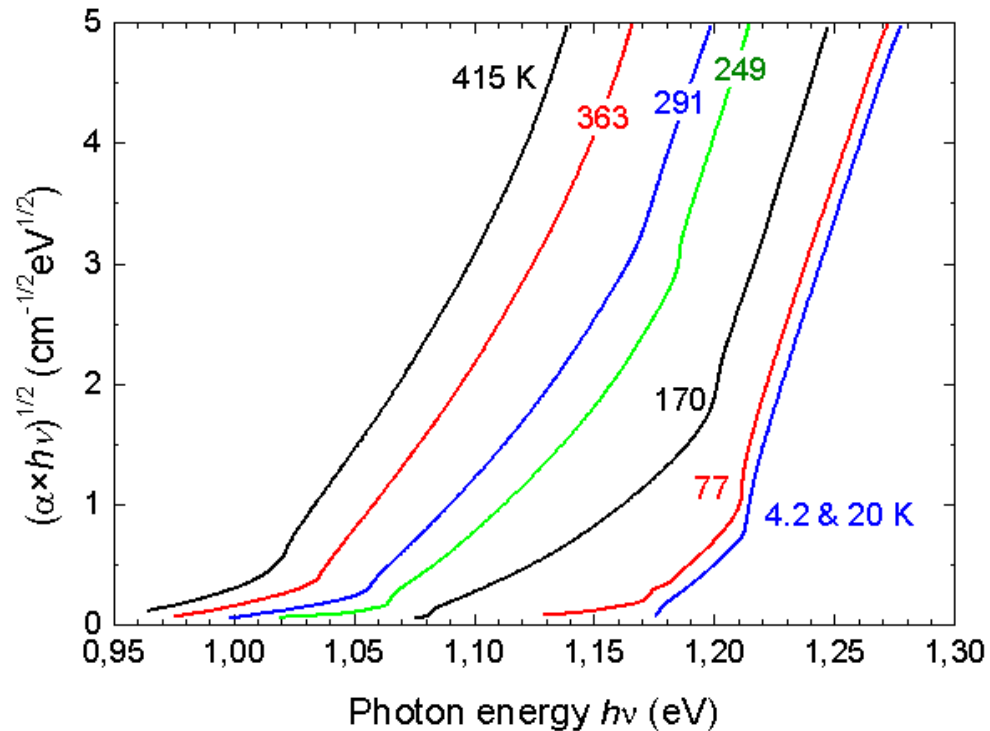
- Wafer with fast surface recombination
- With a long carrier lifetime ( $> 100 \mu\text{s}$ )



... See the diffusion-driven surface recombination?

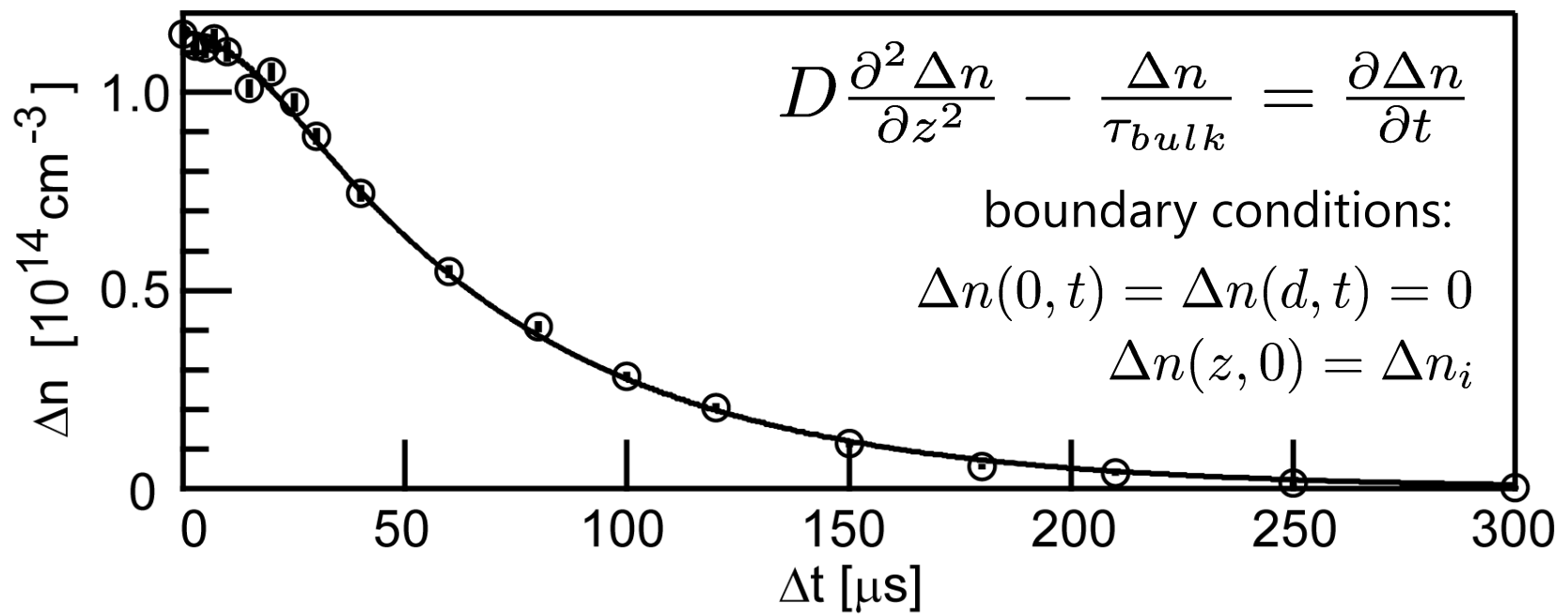
# Uniform carrier injection

To make the analysis simpler, let's choose a low  $\alpha$  for (nearly) uniform carrier injection e.g.  $\alpha = 10 \text{ cm}^{-1}$  gives 1 mm for absorption length



There are tail states, which allows a low  $\alpha$

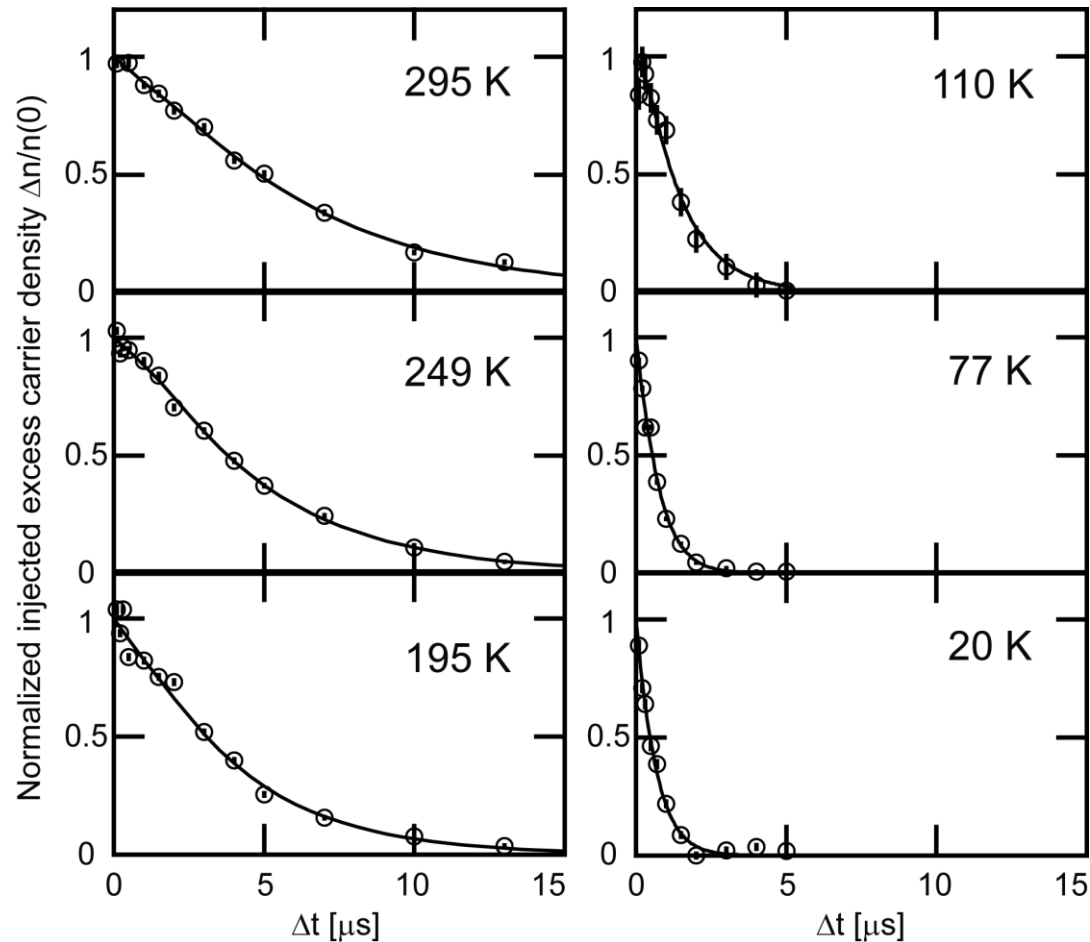
# Excess carrier dynamics



- A shoulder where the surface effect comes in
- Fit with a 1-D diffusion equation for  $\Delta n(z=d/2, t)$  with diffusion constant  $D$  & bulk carrier lifetime  $\tau_{bulk}$  as fit parameters and boundary conditions
- $D$  is the "effective" diffusivity including both  $e$  and  $h$  i.e. ambipolar diffusion

# Temperature dependent measurements

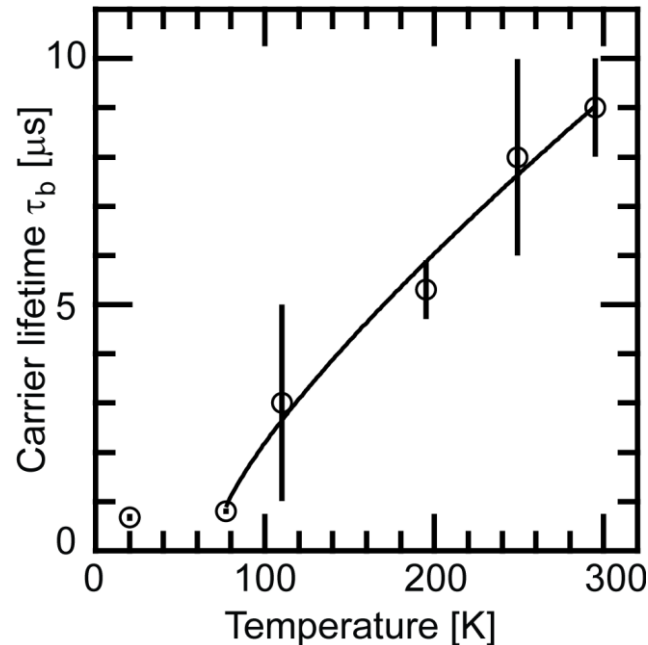
- The method works in Ge in the same way
- Sample was cooled down
- Temperature dependent measurements to get  $\tau_{\text{bulk}}(T)$  & mobility  $\mu(T)$





# Temperature dependent measurements

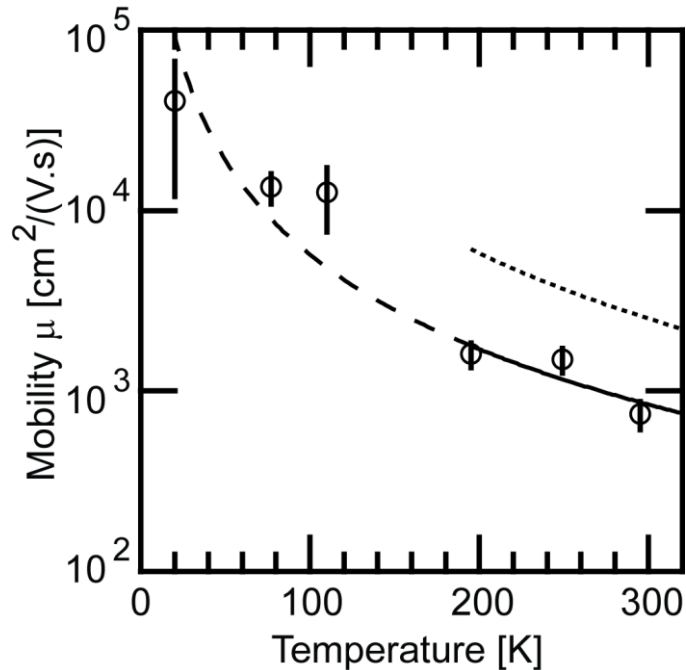
## Bulk carrier lifetime $\tau_{\text{bulk}}$



- $\tau_{\text{bulk}}(T)$  may show recombination mechanisms (& type of defect centres)
- Monotonic decrease with cooling, then stays constant
  - "excitonic Auger recombination"
- A carrier in a free exciton is captured & recombines in an impurity site
- Free exciton formation becomes the rate-limiting step
- All  $e-h$  forms free excitons <76 K

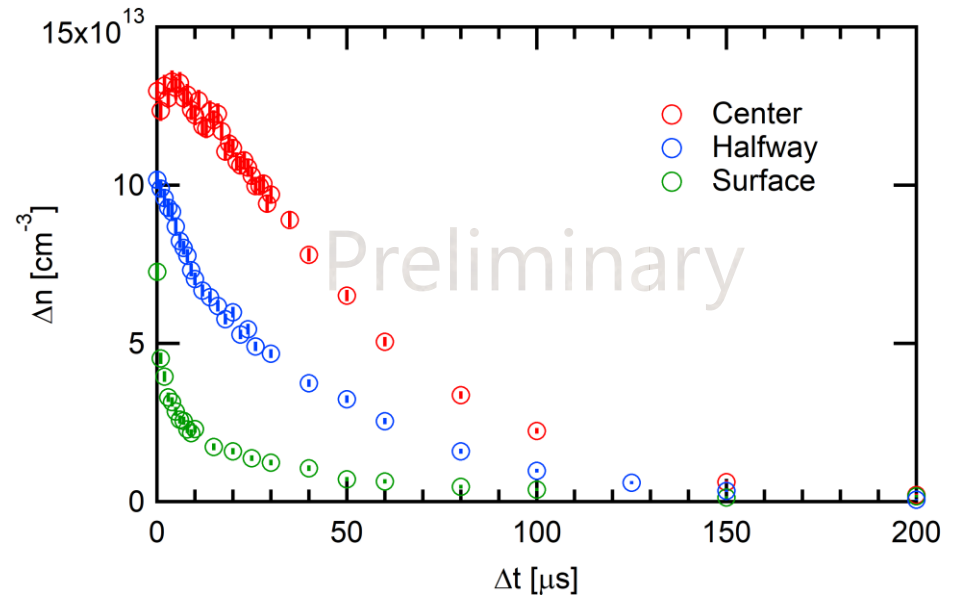
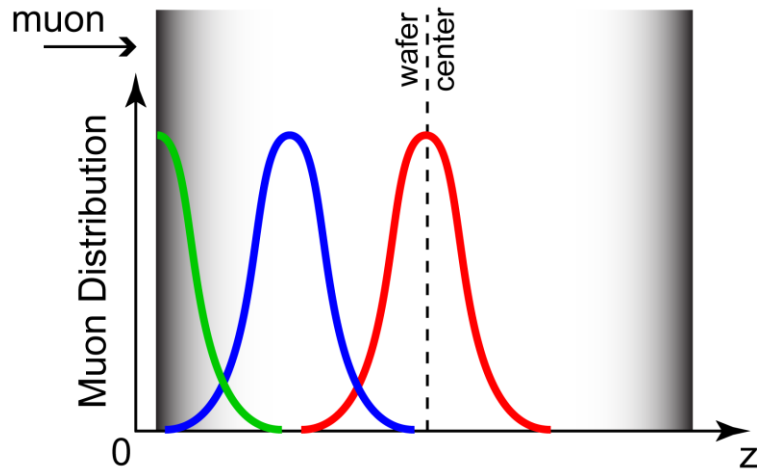
# Temperature dependent measurements

## Carrier mobility $\mu$



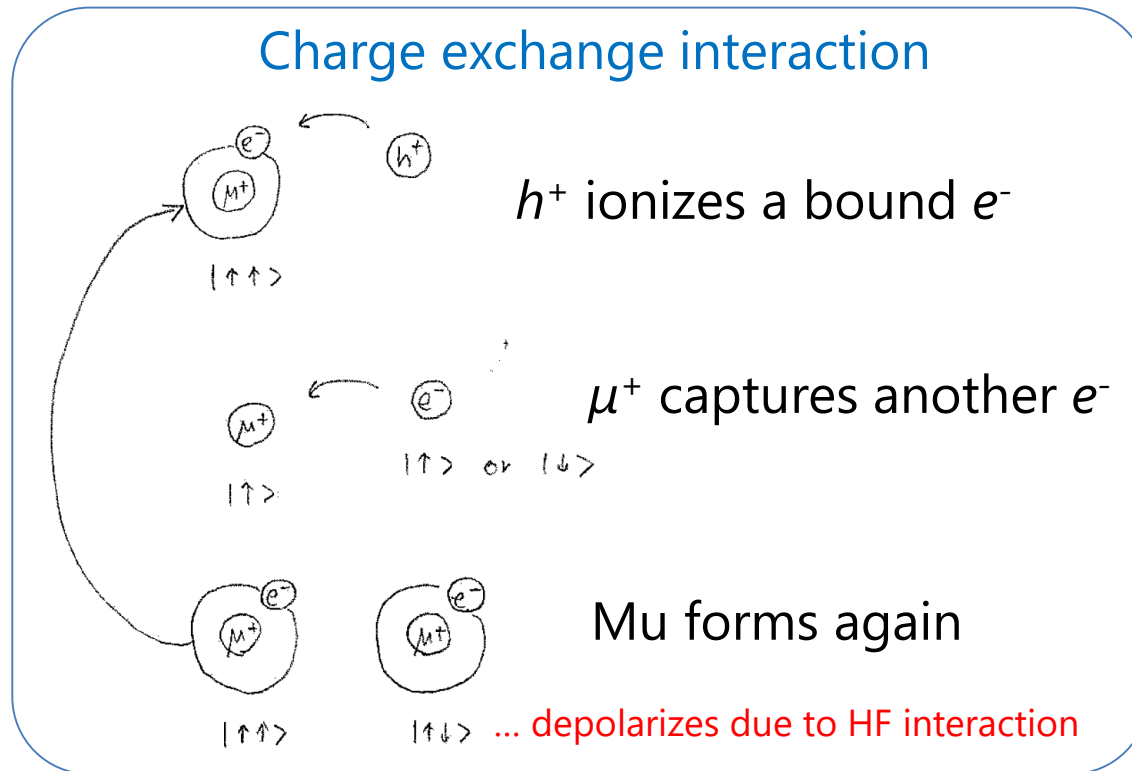
- Monotonic increase with cooling ... less lattice scattering
- In high T, electrons ( $\mu_e$ ) and holes ( $\mu_h$ ) diffuse together
- In low T, exciton diffusion & impurity scattering are dominant
- In 20 K,  $D_{ex} = 69 \text{ cm}^2/\text{s}$  ... comparable with  $D_{ex} = 63 \text{ cm}^2/\text{s}$  from a literature

# Space-resolved measurement



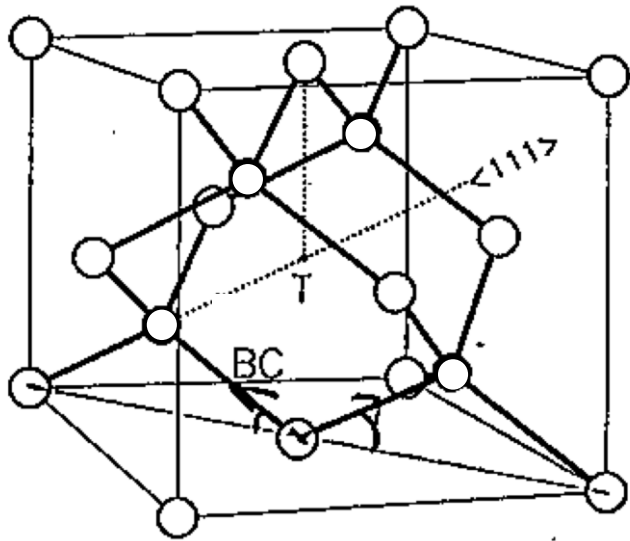
- Silicon, no surface passivation
- Degraders can change the implantation depth  
... can predict the position
- Faster recombination seen closer to the surface
- Spatial resolution is not so great; 130  $\mu\text{m}$  FWHM
- Should help determining the fit parameters more accurately

# Complex Mu dynamics



- This is a oversimplified model
- In the real world it's more complicated ...

# Muonium states in intrinsic silicon



Two Mu sites:

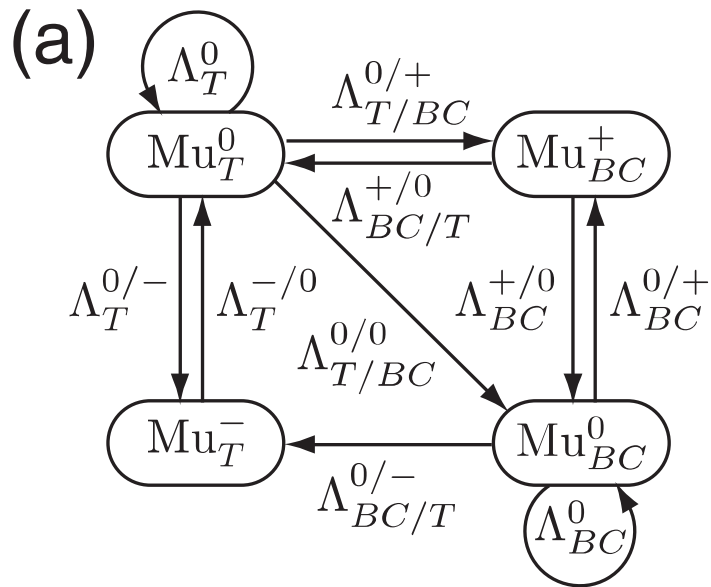
- Tetrahedral interstitial site (T-site)
- Bond-centered site (BC-site)

BC-site:  $\text{Mu}_{\text{BC}}^0$  &  $\text{Mu}_{\text{BC}}^+$

T-site:  $\text{Mu}_{\text{T}}^0$  &  $\text{Mu}_{\text{T}}^-$

They may exchange states dynamically

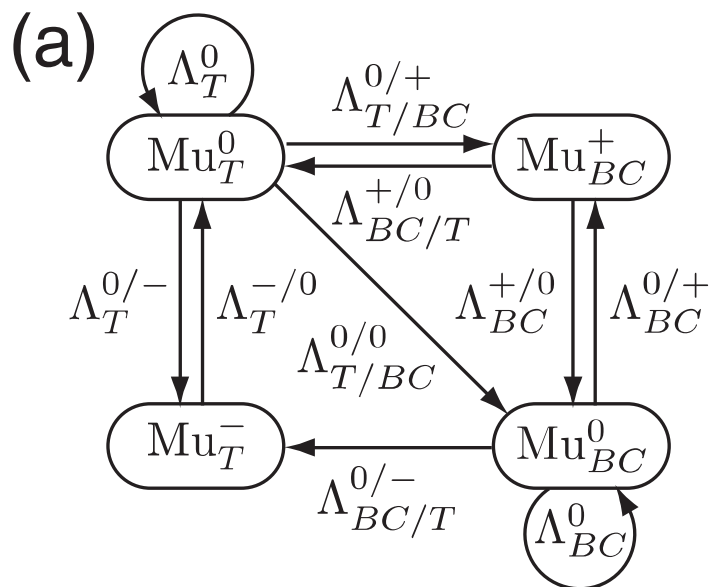
# Carriers activate the network



Complex system...

- Site-change process
- Carrier exchange process
- Spin exchange process

# Carriers activate the network



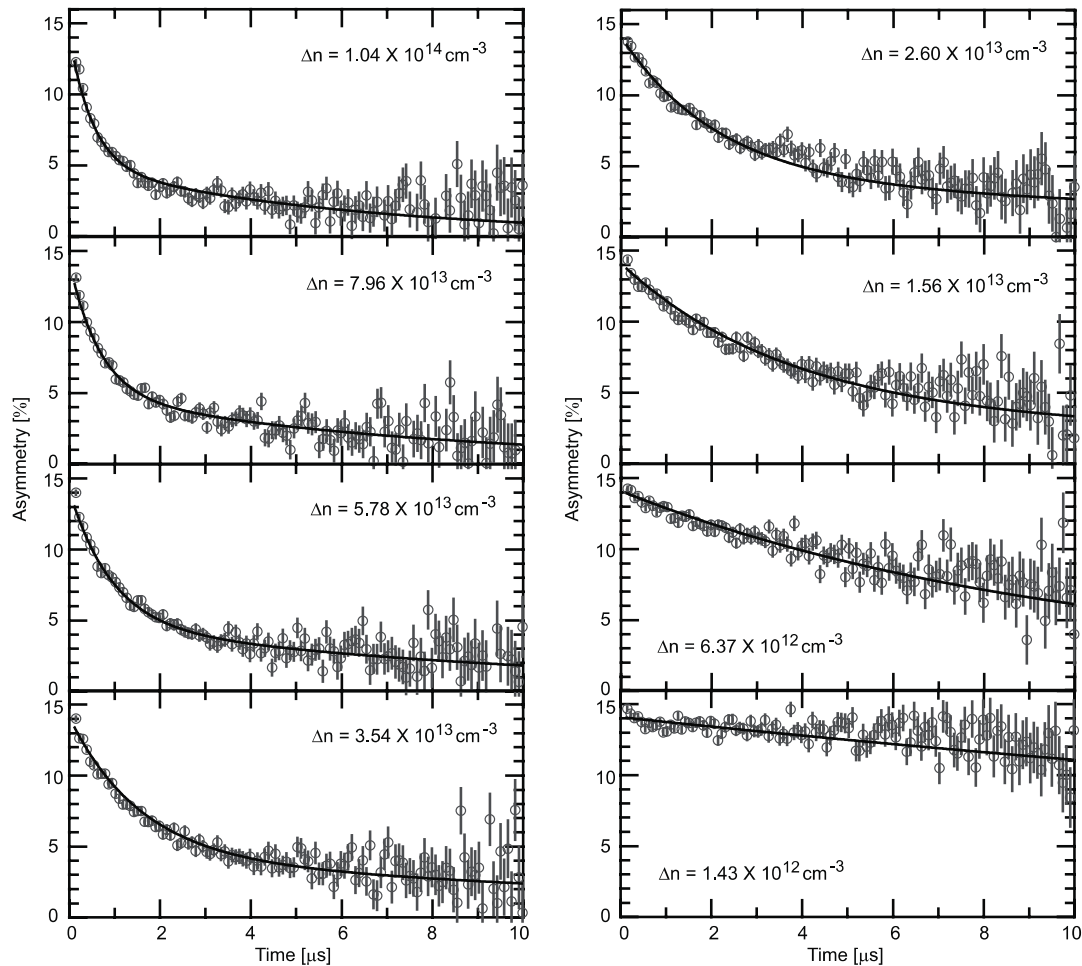
Mu <sub>BC</sub> <sup>0</sup> relaxation	$\Lambda_{BC}^0 = nv_n \sigma_{BC}^0$ $\Lambda_{BC}^{0/+} = \alpha_{BC}^{0/+} \exp[-E_{BC}^{0/+}/k_B T] + pv_p \sigma_{BC}^{0/+}$ $\Lambda_{BC/T}^{0/-} = nv_n \sigma_{BC/T}^{0/-} \exp[-E_{BC/T}^{0/-}/k_B T]$
Mu <sub>BC</sub> <sup>+</sup> relaxation	$\Lambda_{BC}^{+/0} = nv_n \sigma_{BC}^{+/0}$ $\Lambda_{BC/T}^{+/0} = nv_n \sigma_{BC/T}^{+/0} \exp[-E_{BC/T}^{+/0}/k_B T]$
Mu <sub>T</sub> <sup>0</sup> relaxation	$\Lambda_T^0 = nv_n \sigma_T^0$ $\Lambda_{T/BC}^{0/0} = \alpha_{T/BC}^{0/0} \exp[-E_{T/BC}^{0/0}/k_B T]$ $\Lambda_{T/BC}^{0/+} = pv_p \sigma_{T/BC}^{0/+} \exp[-E_{T/BC}^{0/+}/k_B T]$ $\Lambda_T^{0/-} = nv_n \sigma_T^{0/-}$
Mu <sub>T</sub> <sup>-</sup> relaxation	$\Lambda_T^{-/0} = \alpha_T^{-/0} \exp[-E_T^{-/0}/k_B T] + pv_p \sigma_T^{-/0}$

Transitions are characterised by

- Carrier density
- Absorption cross section
- Activation energy
- etc.

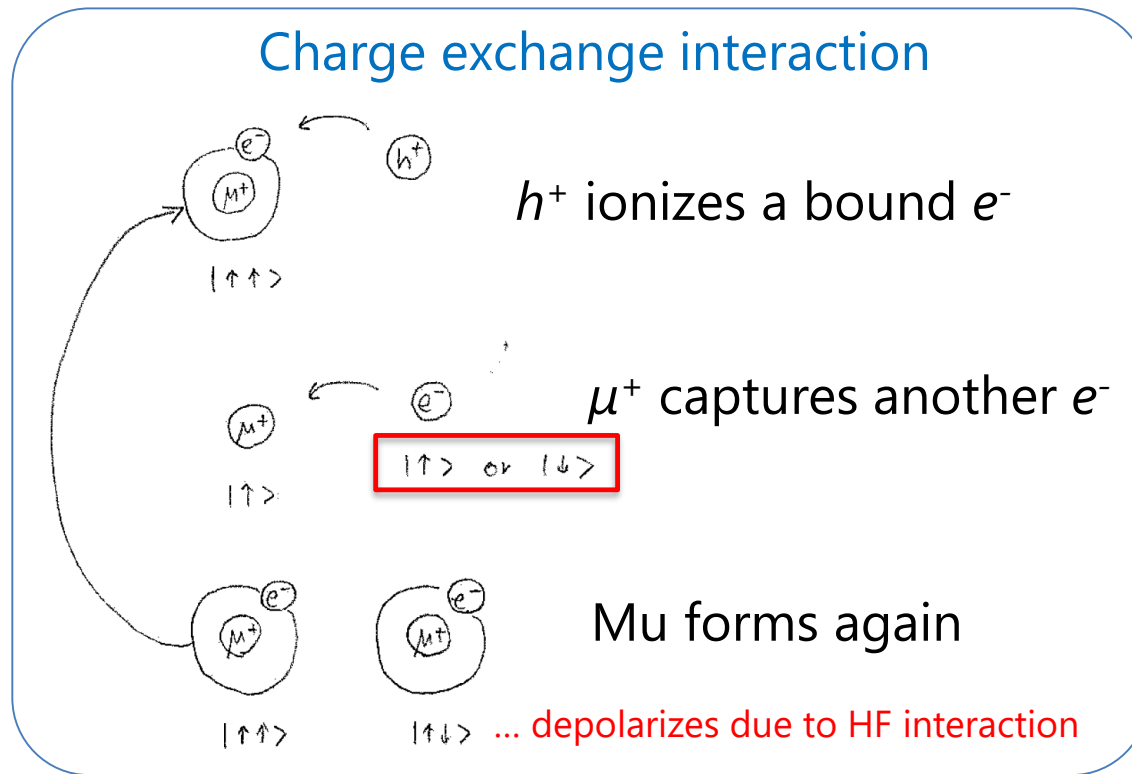
# QUANTUM can solve the mechanism

Simultaneous fit for different carrier densities





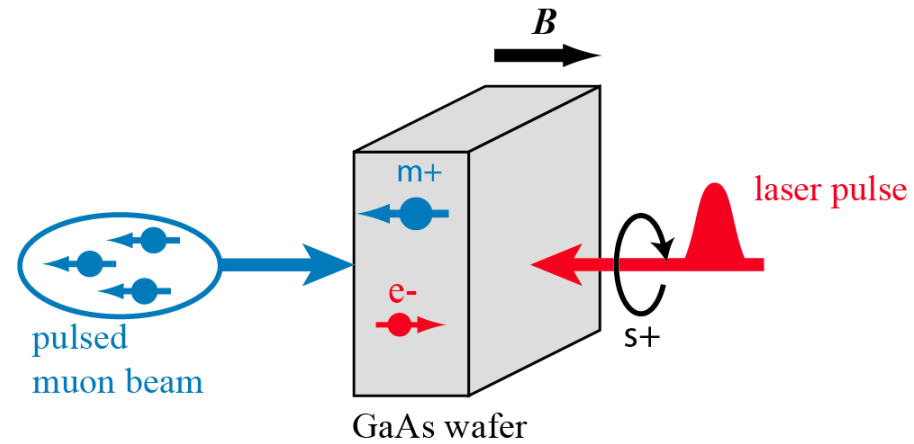
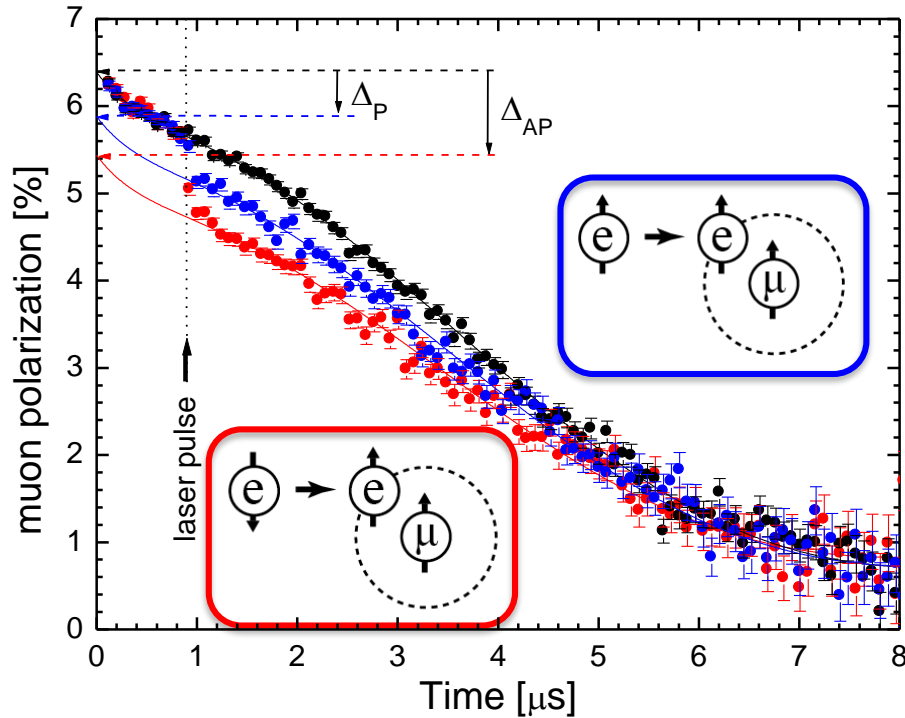
# What if there's more spin-up (or down) $e^-$ ?



- “Optical spin orientation” can choose which spin to excite
  - ... circularly polarized light
  - ... spin angular momentum of light transferred to electronic spin via spin-orbit interaction

# Spin polarisation effect seen in GaAs

## Optical spin orientation & ZF- $\mu$ SR

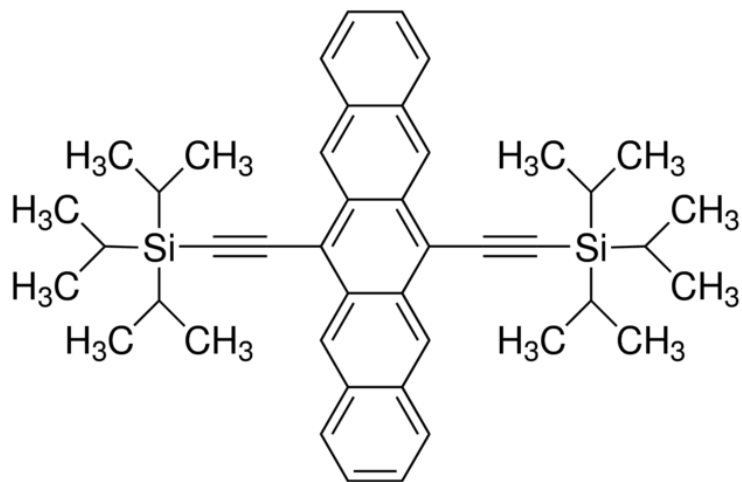


Parallel: less depolarisation

Anti-parallel: more

## Case study #2

### Photoexcitation of TIPS-Pentacene



- Organic semiconductor
- good optoelectronic properties
- Photovoltaic applications?

K. Wang, A.J. Drew, et al. Nat Mater **16**, 467 (2017).

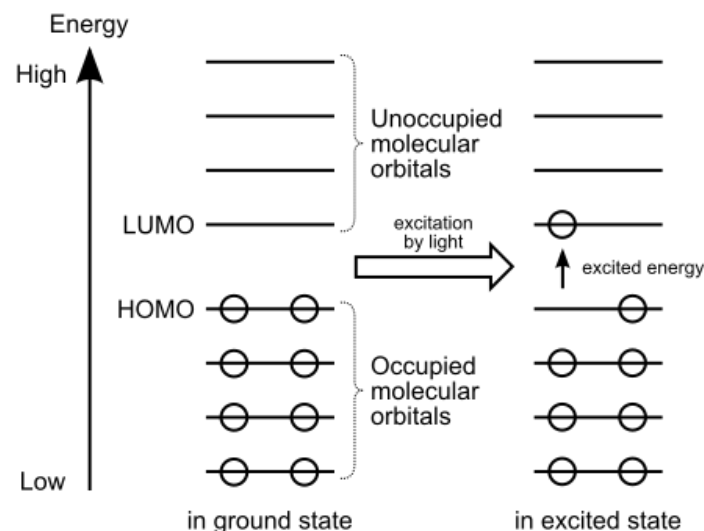
# Photoexcite an organic molecule

Bandgap in organic molecule:

HOMO-LUMO gap

Highest Occupied Molecular Orbit

Lowest Unoccupied Molecular Orbit



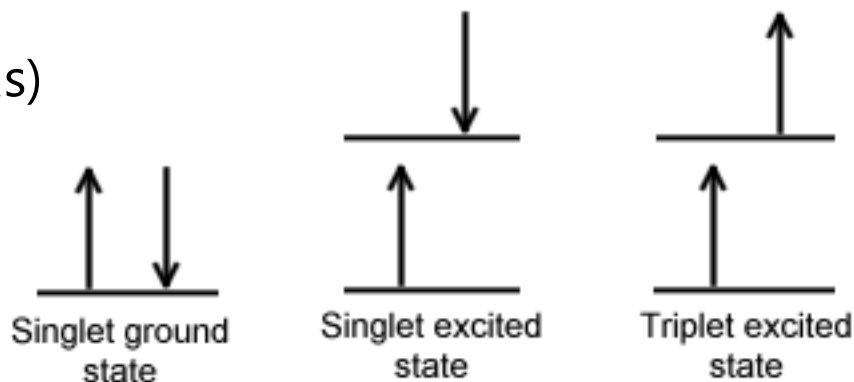
Singlet excited state (short lived,  $\sim$ ns)

transferred to

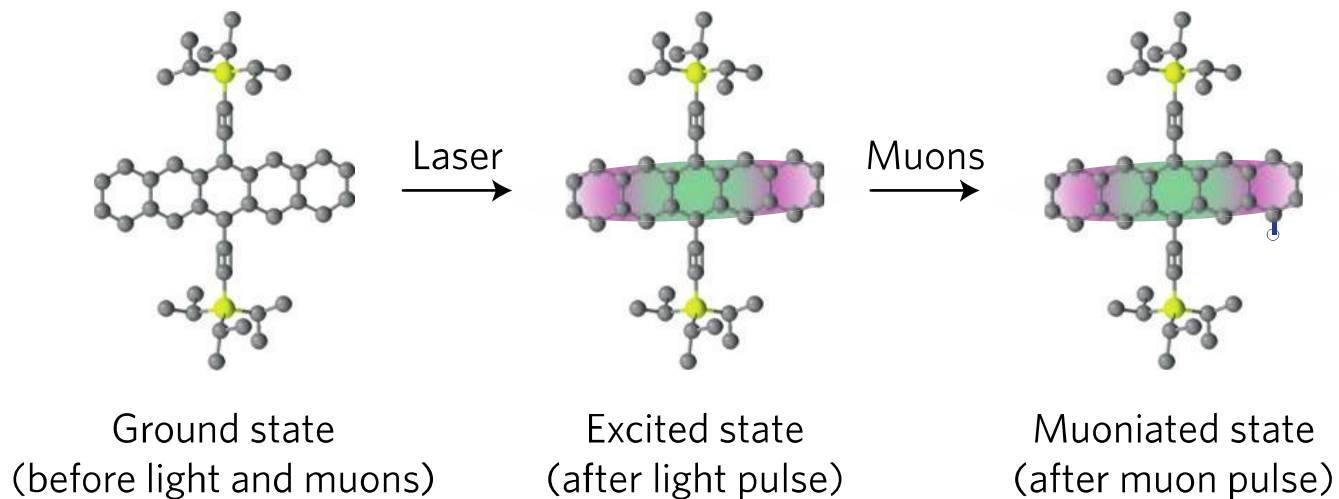
triplet excited state (longer lifetime,  $\sim$  $\mu$ s)

via intersystem crossing

... can we measure the kinetics using muons?



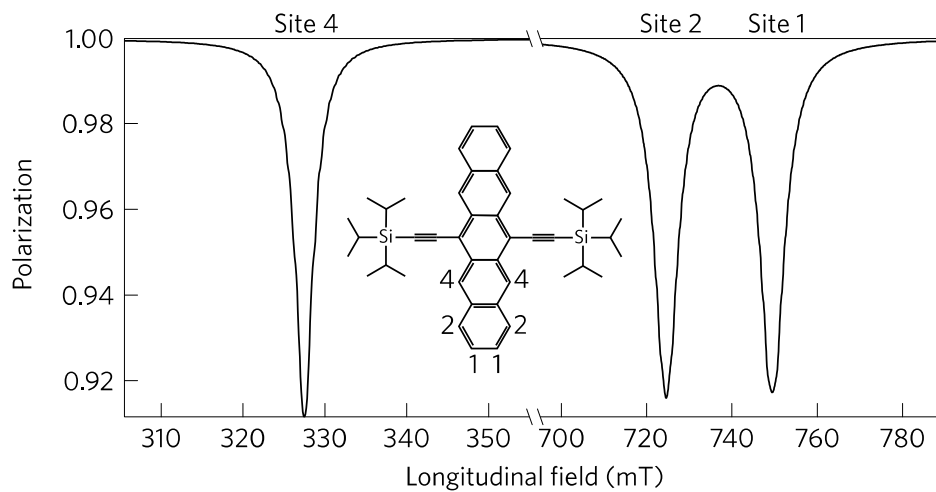
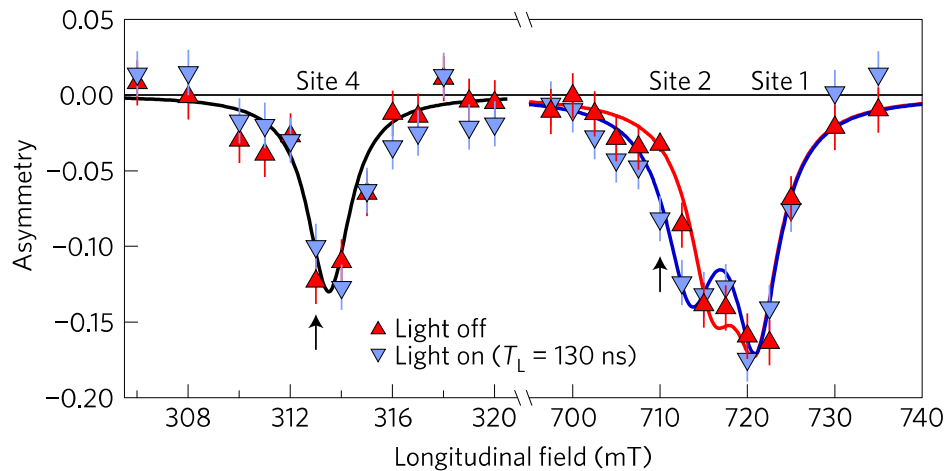
# What does muon see ?



- Excited molecule has different electron wave function
- Should change the ALC signal

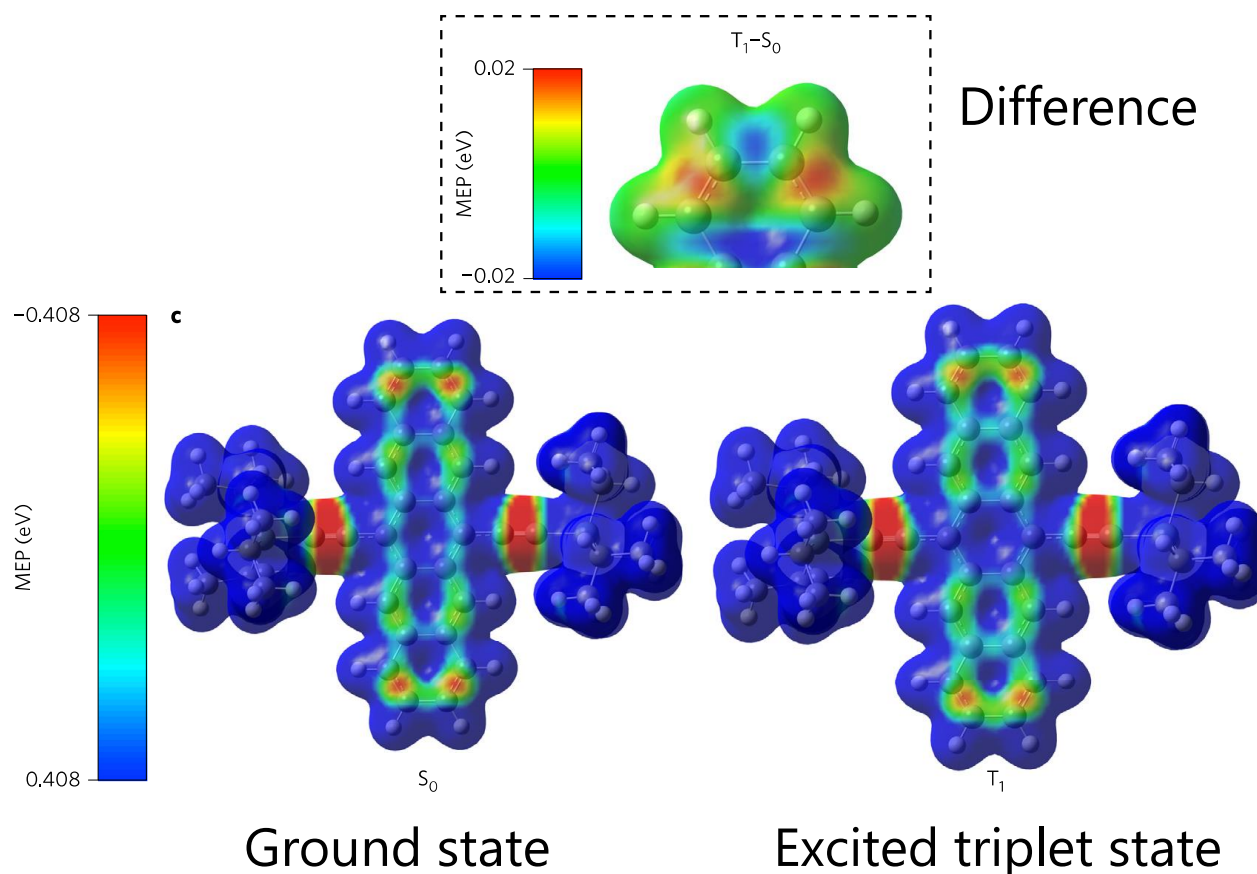
# Change in ALC spectrum

- Site assignment was done using DFT calculation
- Photoinduced change was in Site 2



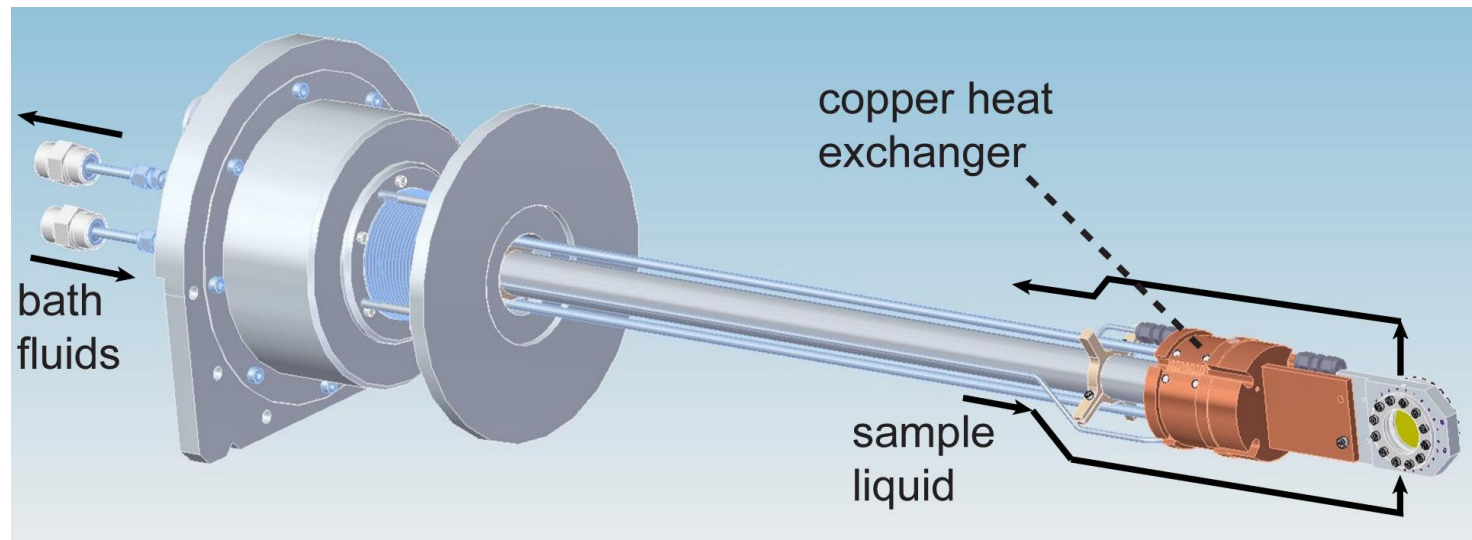
# How can we understand it

- DFT calculation predicts more electron density near Site 2
- Site 2 becomes more reactive



# Difficulties in exciting liquid samples

- Poor heat dissipation ... possible local heating
- Lower diffusivities ... can Mu find excited molecules ?
- More concentration, shorter optical absorption length, less muon/photon overlap



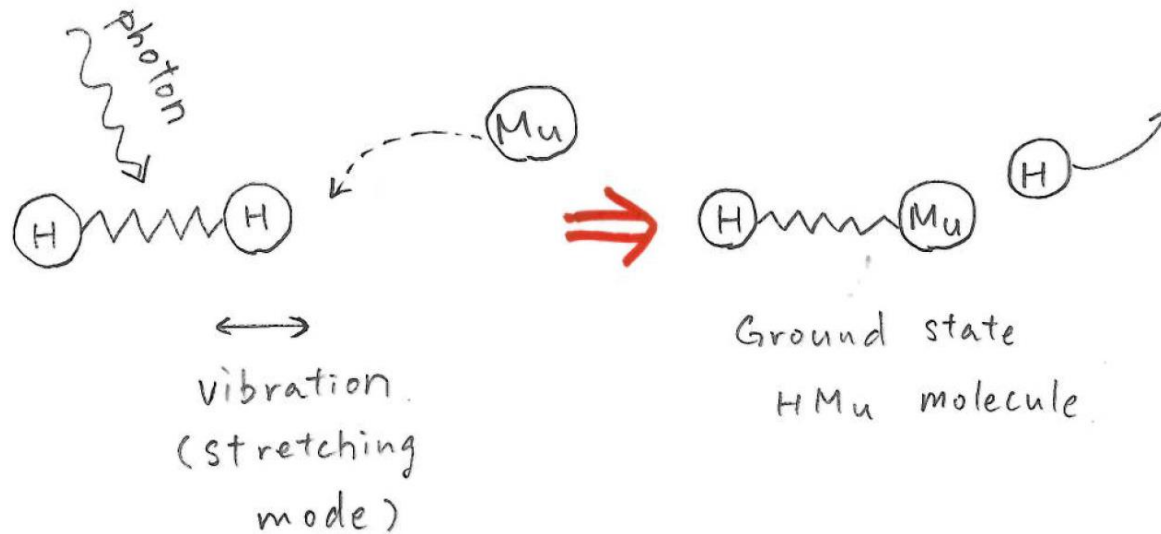


## Case study #3

### Reaction kinetics of $\text{Mu} + \text{H}_2$

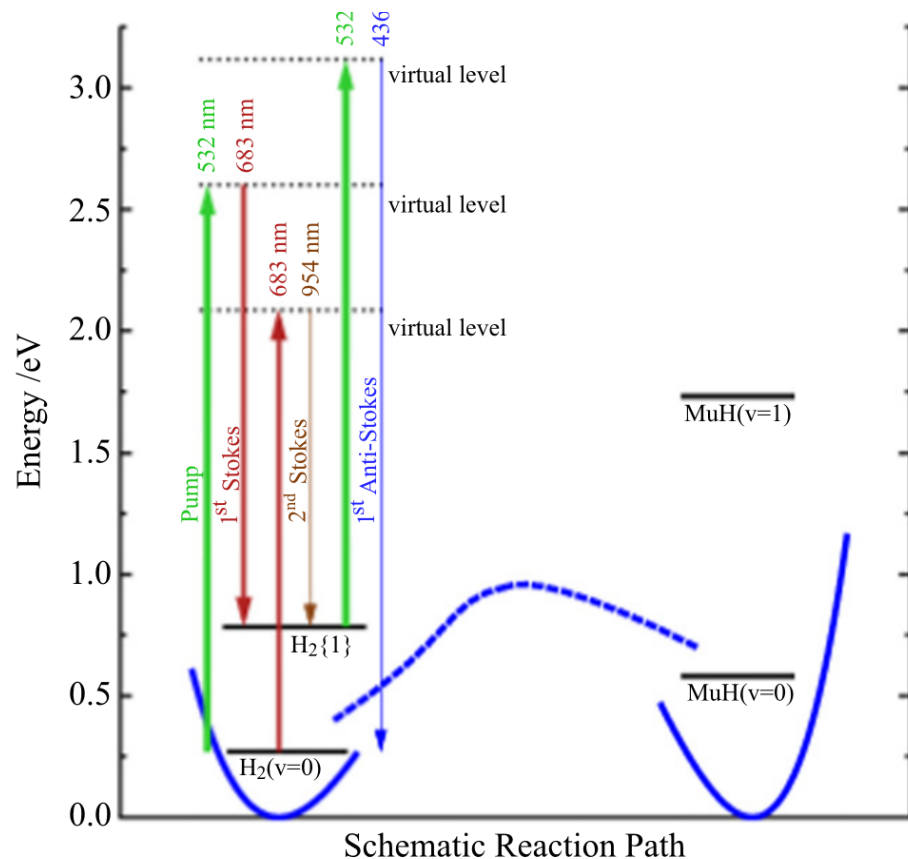
- P. Bakule, D. G. Fleming, et al. J. Phys. Chem. Lett. 3, 2755 (2012).
- P. Bakule, D. Fleming, et al. J. Phys. B: At. Mol. Opt. Phys. 48, 045204 (2015).

# Summary of the experiment



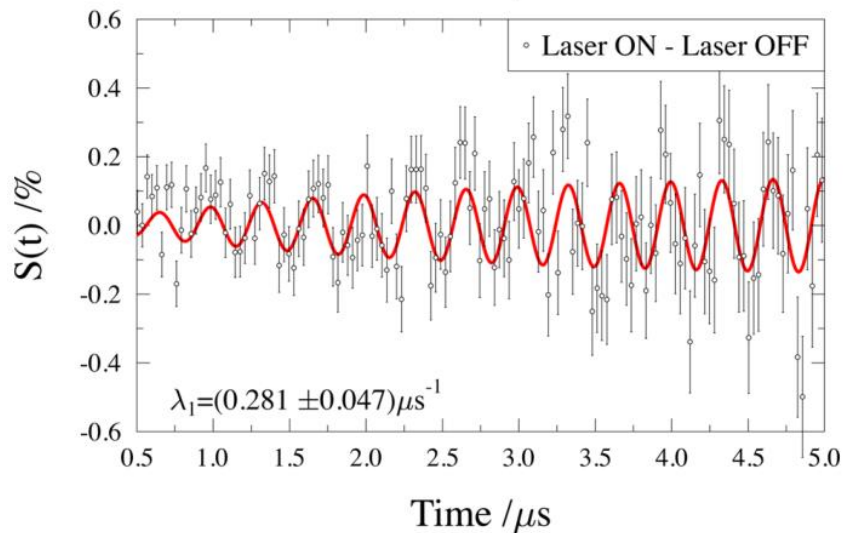
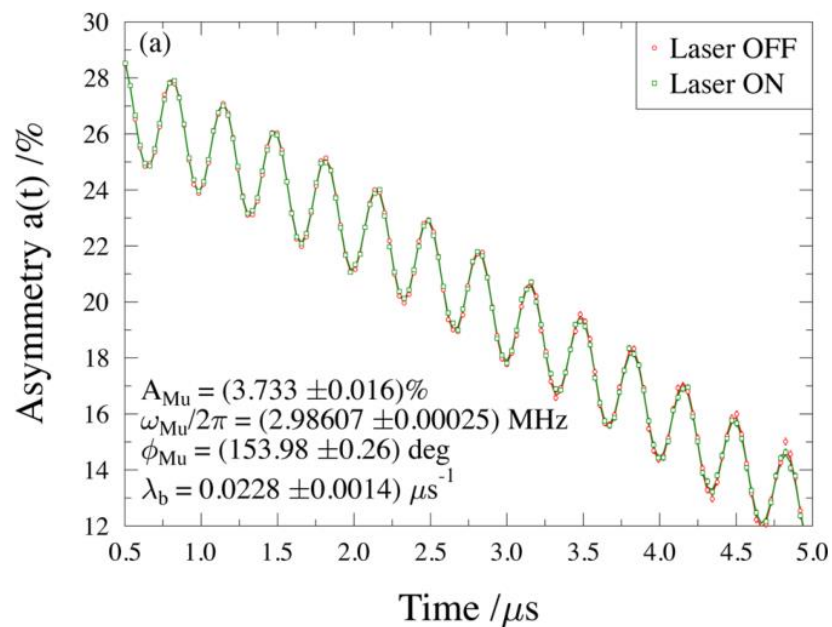
- Photons excite  $\text{H}_2$  molecules to the first vibrational excited state
- Reacts with Mu, forming H-Mu molecules
- Observed in Mu precession signal as reduction of Mu fraction (b/c Mu in H-Mu is diamagnetic)
- Need pressurized gas

# How to excite gaseous molecules?



- Wavelength for  $H_2(v=0) \rightarrow H_2(v=1)$  is  $\sim 2.4 \mu\text{m}$  ... difficult to get a high-power light source
- Simulated Raman Scattering ... generates wavelength different from the pump ... the energy difference is associated with vibrational/rotational/electronic states in material ... used to be a method of wavelength conversion
- Can populate the  $H_2(v=1)$  state

# Summary of the experiment



- Mu precession under TF 2G
- Light arriving at  $\Delta t = 0$
- Laser OFF/ON ... there's a difference
- Laser (ON) – (OFF) makes it more clear
- The difference becomes larger b/c H-Mu population grows
- Can obtain the reaction rate for  
$$\text{Mu} + \text{H}_2\{1\} \xrightarrow{\lambda_1} \text{MuH} + \text{H}$$
- Test bench for the first-principle calculation

# Planning an experiment

But you have to carefully plan your experiment:

- Muon/Photon spatial overlap
- How do photon/muon interact?
  - Sometimes you have to fight diluted nature of muon
  - e.g. photocarriers in Si diffuse and “find” Mu
- Ideal to know either optical character or muon signal
- More heating in lower T (heat capacity proportional to  $T^3$ )

Photons are everywhere ... idea ??