

Multigap Superconductivity in Chiral Noncentrosymmetric TaRh_2B_2

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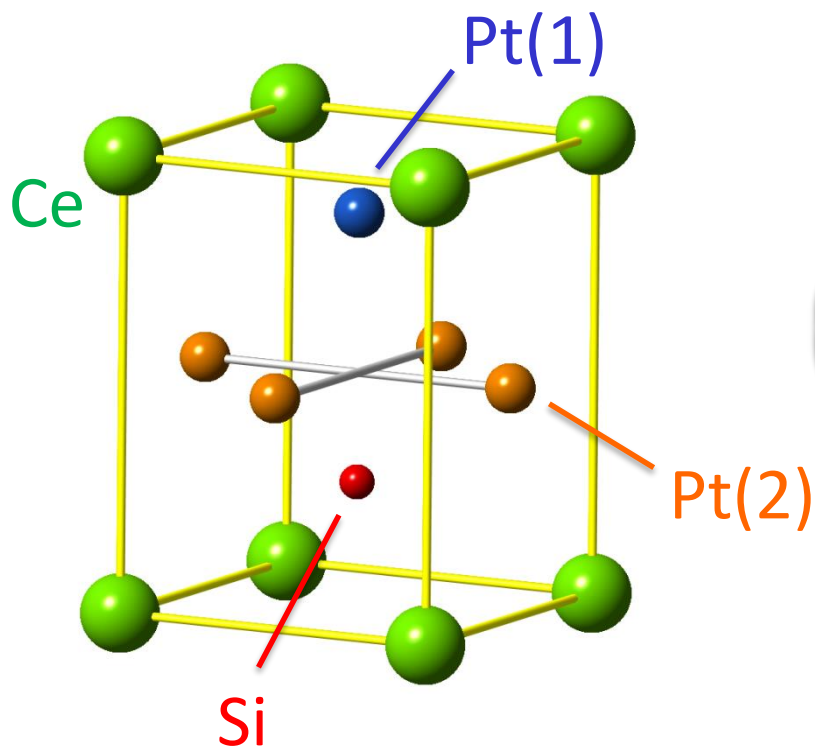


Inversion Symmetry

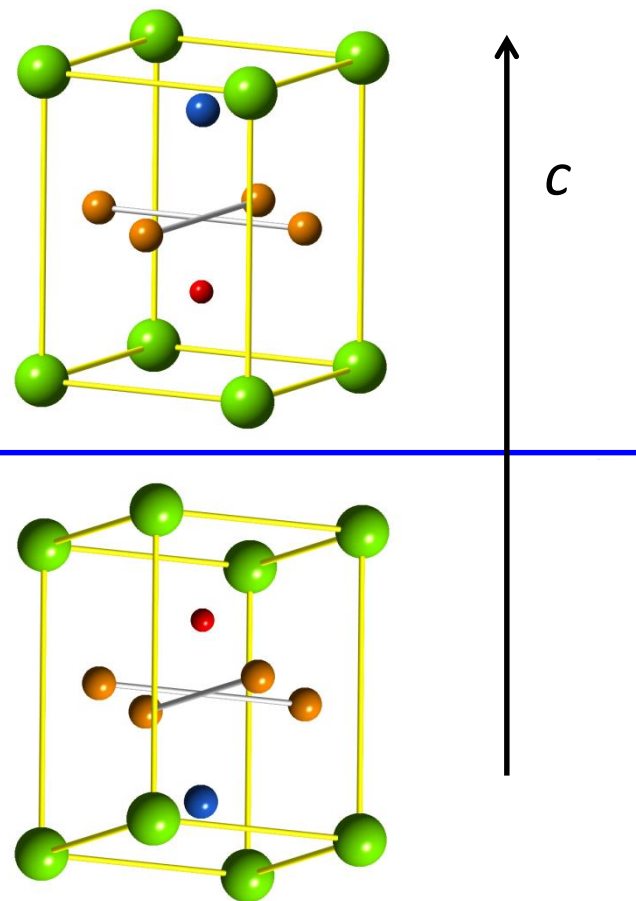
(or lack thereof...)

$(x, y, z) \rightarrow (-x, -y, -z)$ not a symmetry of the crystal.

CePt₃Si



ab



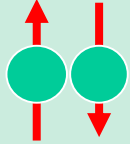
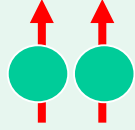
c

Superconductivity

The gap function of a superconductor has the form

$$\hat{\Delta}(\mathbf{k}) = \underbrace{\begin{bmatrix} 0 & \Delta_0 \\ -\Delta_0 & 0 \end{bmatrix}}_{\text{Singlet}} + \underbrace{\begin{bmatrix} -d_x + id_y & d_z \\ d_z & d_x + id_y \end{bmatrix}}_{\text{Triplet}}$$

If the crystal has inversion symmetry then we have:

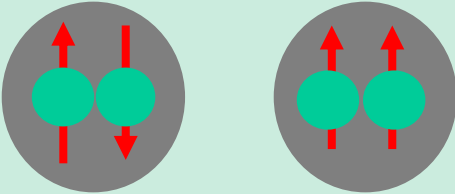
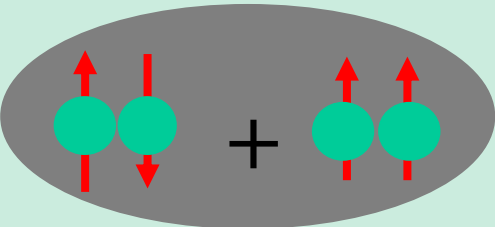
	Spin part	Orbital part
Singlet	antisymmetric 	Symmetric (s, d...)
Triplet	symmetric 	Antisymmetric (p, f...)

But what if the crystal lacks a centre of inversion?

Noncentrosymmetric Superconductivity

Rashba-type antisymmetric spin-orbit coupling (ASOC) allows degeneracy of the Fermi surfaces.

Parity no longer a good quantum number.

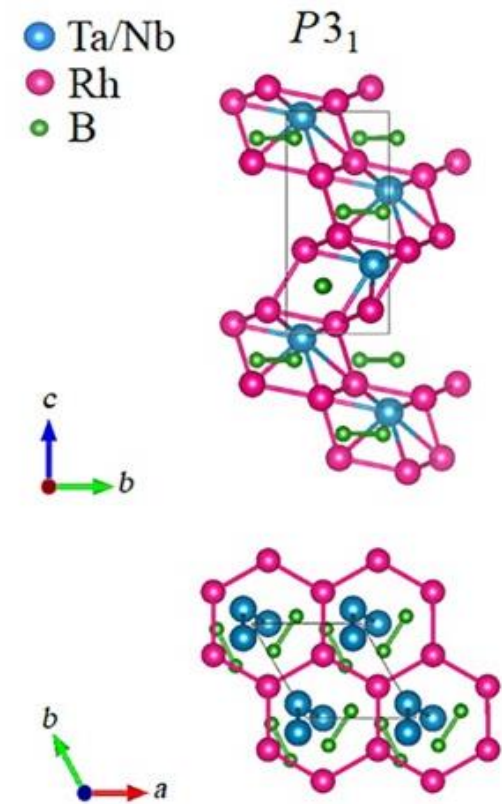
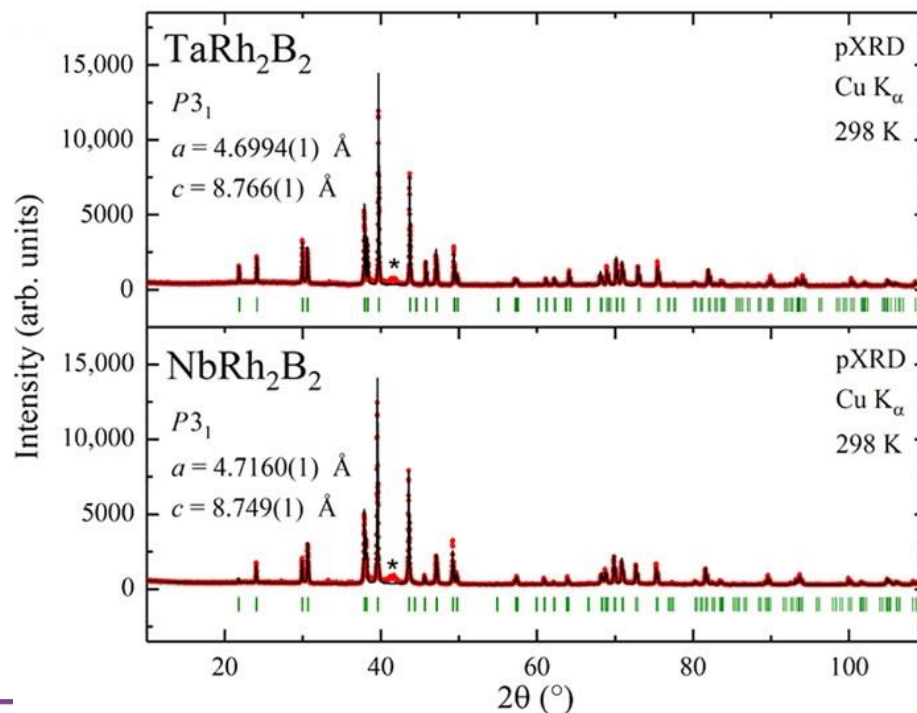
	Centrosymmetric	Noncentrosymmetric
Cooper pair		

Cooper-pair wave function may now be a **MIXTURE** of spin-singlet and spin-triplet pairing.

TaRh₂B₂ and NbRh₂B₂

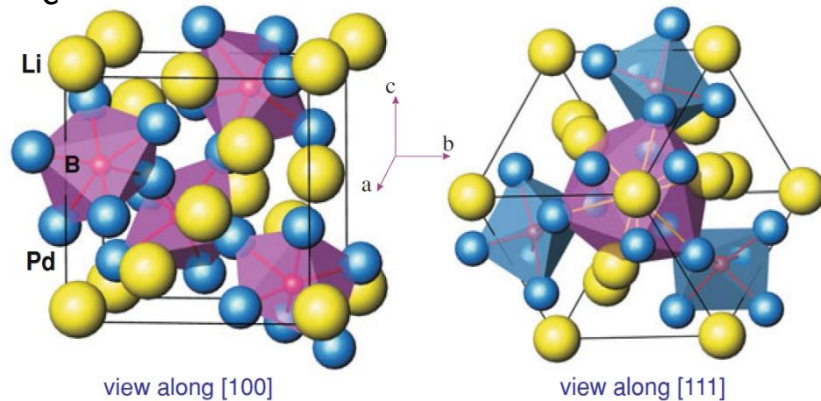
TaRh₂B₂ and NbRh₂B₂ crystallizes in chiral noncentrosymmetric crystal structure.

- Space group: $P3_1$
- Polar & Asymmetric



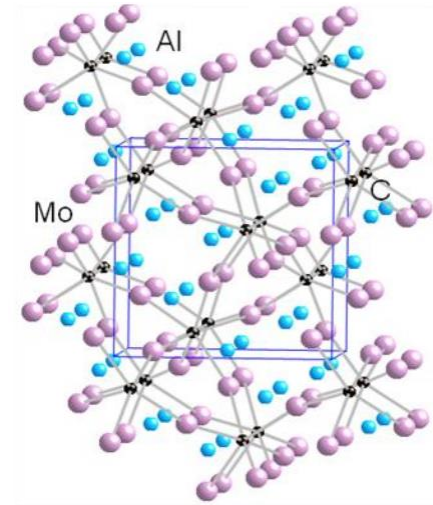
Chirality in Superconductors

$\text{Li}_2\text{Pd}_{1-x}\text{Pt}_x\text{B}$
 $T_c \sim 2.7 - 7 \text{ K}$



Space group: $P4_332$.
 s -wave (Pd) to line nodal gap (Pt).
H. Q. Yaun *et al.* Phys. Rev. Lett. **97**, 017006 (2006).

$\text{Mo}_3\text{Al}_2\text{C}$
 $T_c \sim 9 \text{ K}$



Space group: $P4_132$.
Nodal gap.
A. B. Karki *et al.* Phys. Rev. B **82**, 064512 (2010).

Both noncentrosymmetric with chiral crystal structure.

Experimental Methods

Polycrystalline samples prepared by pressing mixed powder into a pellet and heating in high vacuum furnace at 1200 °C for 6 hours.

Powder XRD confirm:

- structure.
- no impurities including oxides.

Magnetization, heat-capacity, and resistivity.

Muon spectroscopy.

Muon Spin Relaxation/Rotation (μ SR)

Zero-field, longitudinal-field and transverse-field studies were carried out using MuSR spectrometer at ISIS-RAL.

Powder samples mounted on a Ag plate.

^3He cryostat $0.3 \leq T \leq 10$ K.

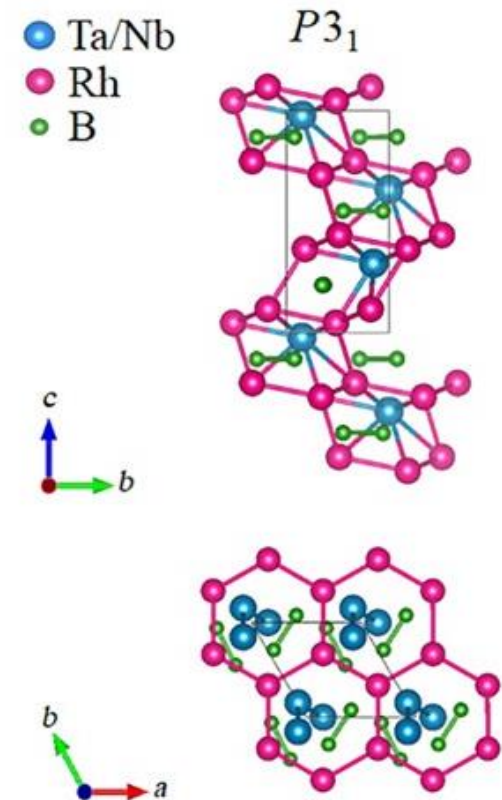
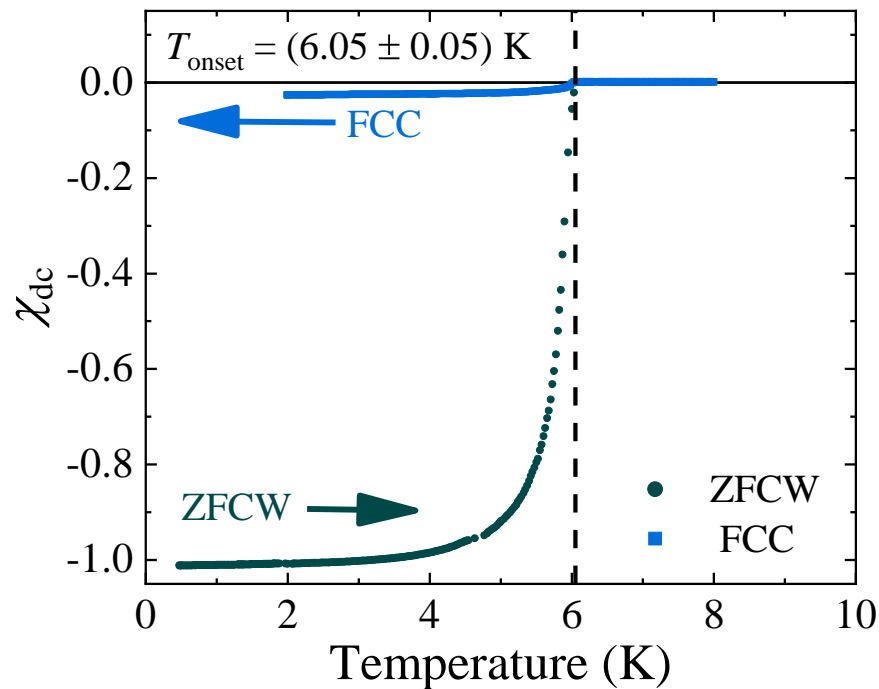
Stray fields cancelled to level of $1 \mu\text{T}$.



TaRh₂B₂

TaRh₂B₂ crystallizes in a chiral noncentrosymmetric crystal structure.

- Space group: $P3_1$
- Polar & Asymmetric

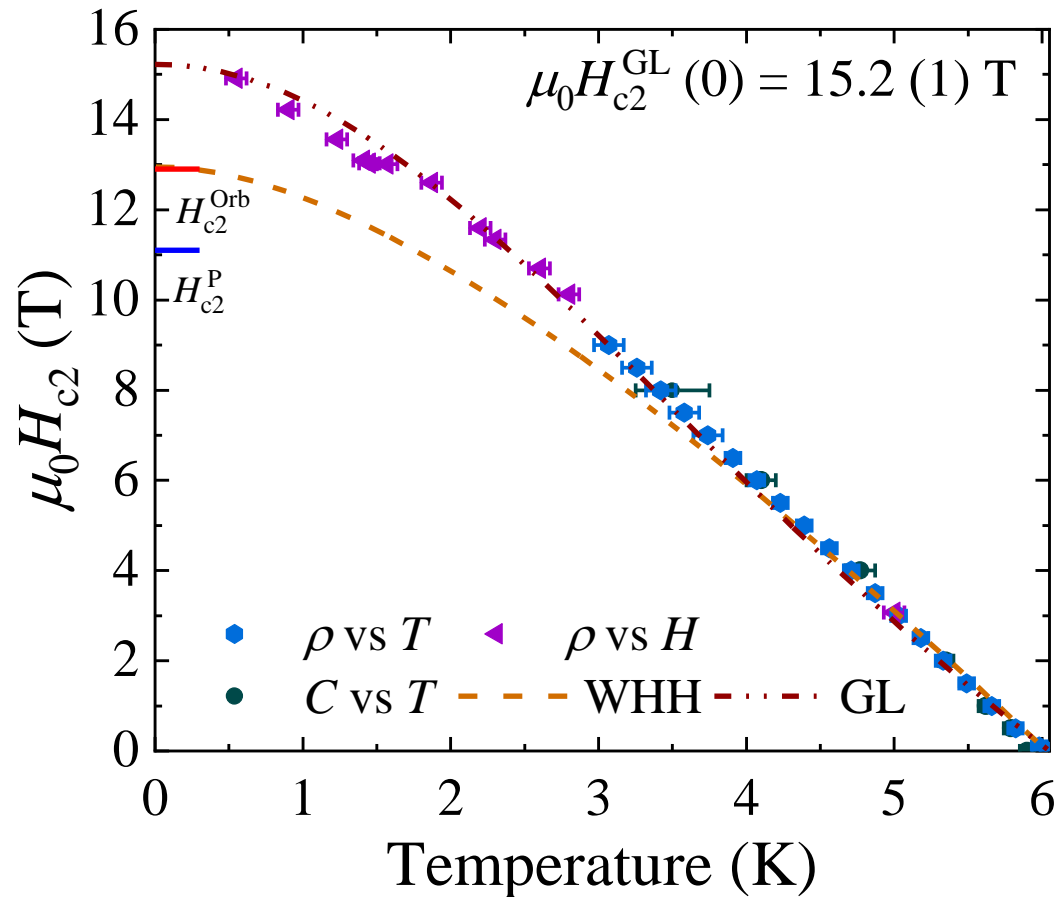


Upper Critical Field

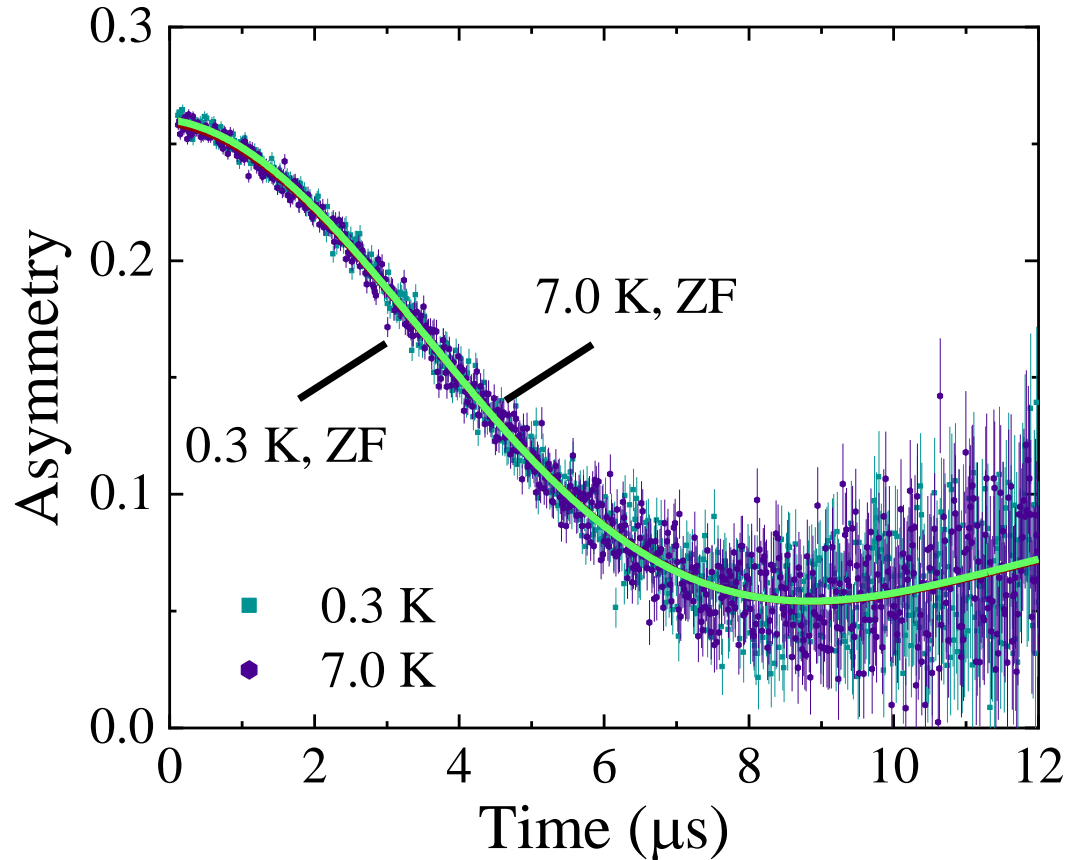
$$\mu_0 H^P = 1.85 \times T_c$$

$$\mu_0 H^P = 11.1 (1) \text{ T}$$

Pauli Limit Violation



Zero-field μ SR



No time-reversal symmetry breaking.

Transverse-field μ SR

TF-spectra in 30 mT.

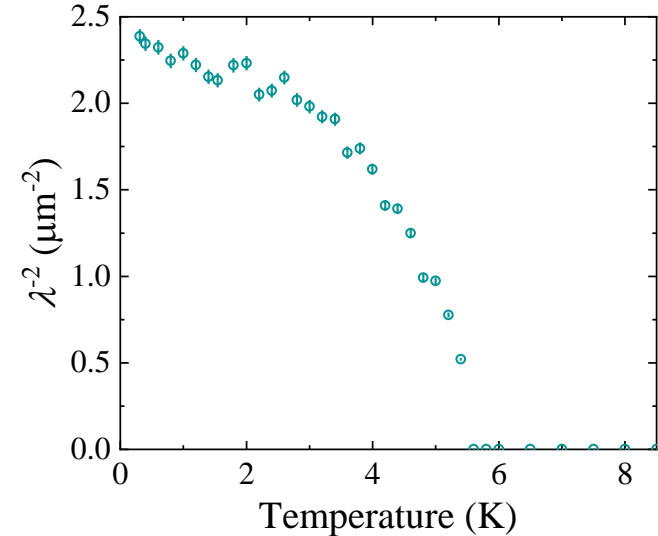
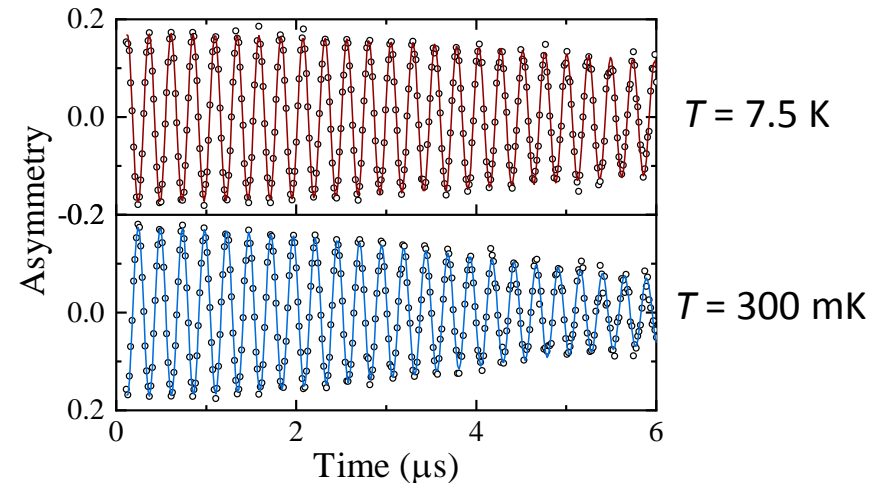


Muons implanted in the silver sample holder do not depolarize.

Spectra fitted with Gaussian relaxation,

$$G(t) = A_1 \exp\left(\frac{-\sigma_{sc}^2 t^2}{2}\right) \cos(2\pi\nu_1 t + \phi_1) + A_2 \cos(2\pi\nu_2 t + \phi_2).$$

σ_{sc} directly related to superfluid density.

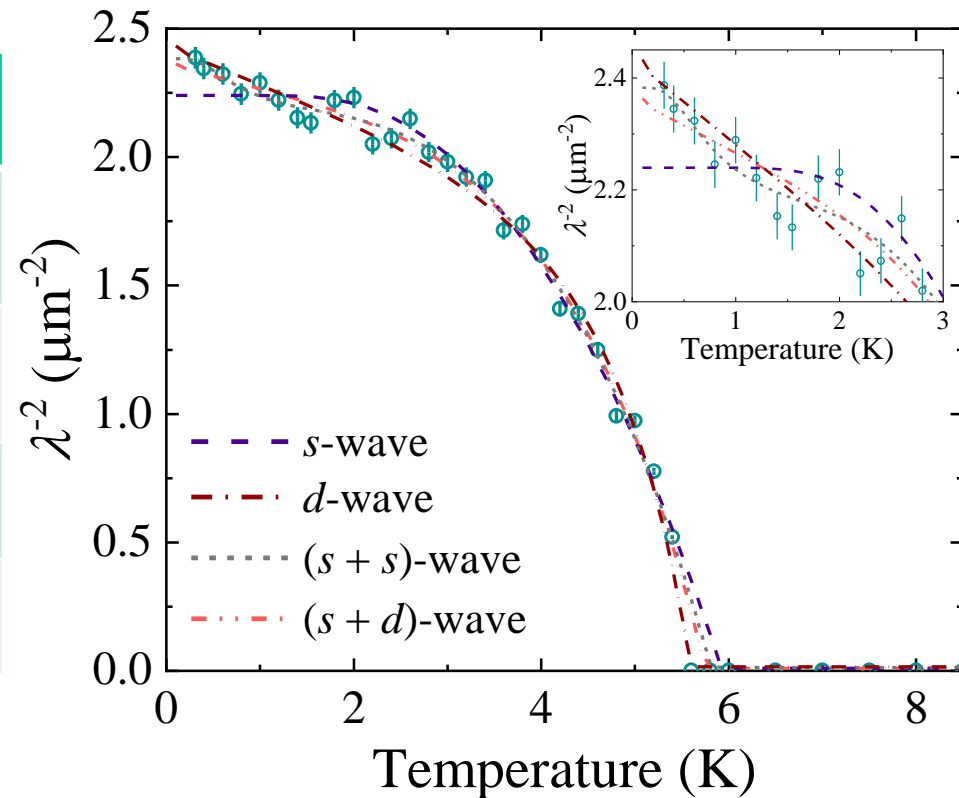


Transverse-field μ SR

$$\left[\frac{\lambda^{-2}(T, \Delta_0)}{\lambda^{-2}(T, \Delta_0)} \right] = w \left[\frac{\lambda^{-2}(T, \Delta_0)}{\lambda^{-2}(T, \Delta_0)} \right] + (1 - w) \left[\frac{\lambda^{-2}(T, \Delta_0)}{\lambda^{-2}(T, \Delta_0)} \right]$$

	<i>d</i> -wave	(<i>s</i> + <i>s</i>)-wave	(<i>s</i> + <i>d</i>)-wave
$\frac{\Delta_{0,1}}{k_B T_c}$	3.78 (14)	2.28 (8)	2.56 (11)
$\frac{\Delta_{0,2}}{k_B T_c}$	---	0.27 (9)	2.83 (9)
<i>w</i>	1	0.89 (2)	0.51 (2)
χ_{norm}^2	1.81	1.18	1.35

cf. BCS value 1.77.



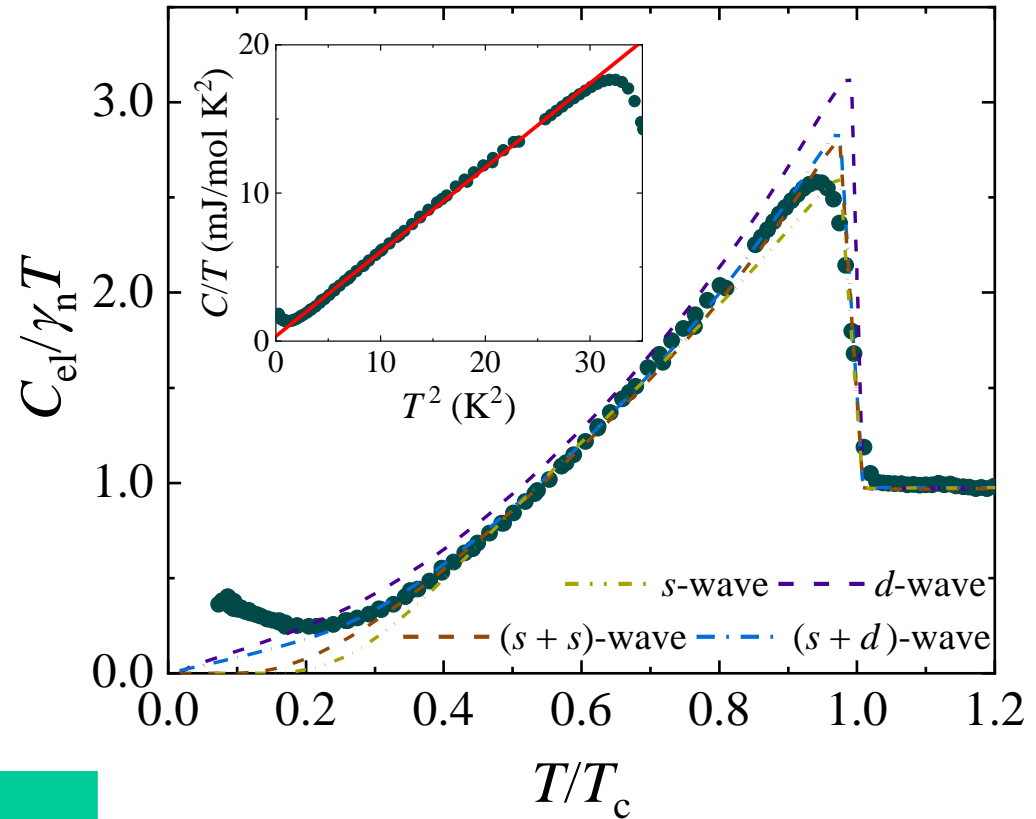
Heat Capacity

Complicated by hyperfine contribution

Clearly not *s*-wave.

(*s+s*)- wave:

$\frac{\Delta_{0,1}}{k_B T_c}$	$\frac{\Delta_{0,2}}{k_B T_c}$	w
2.19 (12)	1.1(3)	4:1



$$\frac{\Delta C}{\gamma_n T_c} = 1.57(2) > 1.43$$

Conclusion

- New chiral noncentrosymmetric superconductor.
- Strong evidence of multigap superconductivity in TaRh_2B_2 .
- Large upper critical field well above Pauli limit.

